



$np \rightarrow d\pi^0$ in chiral perturbation theory —a status report

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for the IOU collaboration

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Charge symmetry breaking

QCD Lagrangian almost symmetric under $u \leftrightarrow d$ exchange
(Charge Symmetry, CS), $P_{CS} = \exp(i\pi\tau_2/2)$
broken by $m_u \neq m_d$, Charge Symmetry Breaking (CSB)
Special case of isospin violation

For hadrons and nuclei CS implies

$$p \rightarrow -n$$

$$n \rightarrow p$$

$$d \leftrightarrow \bar{d}$$

$$\pi^0 \leftrightarrow -\pi^0, \text{ etc.}$$

CSB reviews:

Miller, Nefkens, and Šlaus, PRt194, 1 (1990);

Miller and van Oers, in *Symmetries and Fundamental Interactions in Nuclei*,
Eds. Haxton and Henley, WS, Singapore (1995) p 127

Experimental evidence

- ρ^0 - ω mixing ($e^+e^- \rightarrow \pi^+\pi^-$)
- $n-p$ and other hadron mass differences
- mirror nuclei (e.g. ${}^3\text{He}$ - ${}^3\text{H}$) binding energy, N-S anomaly
- $a_{nn} \neq a_{pp}$
- $np \rightarrow np$: $A_n(\theta_n) \neq A_p(\theta_p)$ analyzing powers
- $dd \rightarrow \alpha\pi^0$ (IUCF)
- $np \rightarrow d\pi^0$ forward-backward asymmetry (TRIUMF)

Neutron-proton mass difference

$m_n - m_p = \delta M + \bar{\delta}M = 1.29 \text{ MeV}$, natural sizes:

$$\delta M \sim \epsilon m_\pi^2 / \Lambda_\chi \sim \text{few MeV} \text{ — strong term}$$

$$\bar{\delta}M \sim \alpha \Lambda_\chi / \pi \sim -2 \text{ MeV} \text{ — EM term}$$

$$\epsilon = (m_d - m_u) / (m_d + m_u) \sim 1/3$$

Models (e.g., Cottingham sum rule) typically $\bar{\delta}M \sim -1 \text{ MeV}$

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χ PT CSB \Rightarrow syms diff for strong and EM:

$$\mathcal{L}_{CSB} = \frac{\delta M}{2} N^\dagger \left(\tau_3 - \frac{\pi_3 \boldsymbol{\tau} \cdot \boldsymbol{\pi}}{2f_\pi^2} \right) N + \frac{\bar{\delta}M}{2} N^\dagger \left(\tau_3 + \frac{\pi_3 \boldsymbol{\tau} \cdot \boldsymbol{\pi} - \boldsymbol{\pi}^2 \tau_3}{2f_\pi^2} \right) N$$

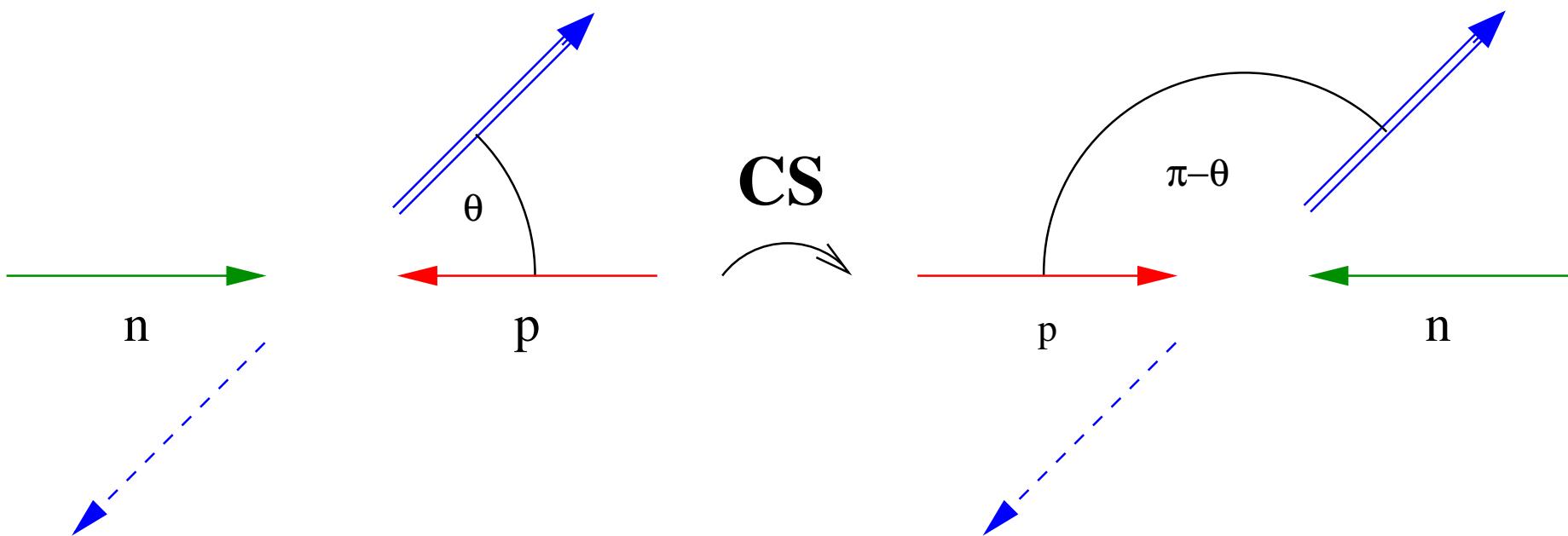
different comb for δM and $\bar{\delta}M$ for $m_n - m_p$ and pions \Rightarrow

Pion reactions to extract $\bar{\delta}M$ and check \mathcal{L}_{CSB} !



CSB in $np \rightarrow d\pi^0$ I

CS $\Rightarrow n \leftrightarrow p$:



$$A_{\text{fb}}(\theta) \equiv \frac{\sigma(\theta) - \sigma(\pi - \theta)}{\sigma(\theta) + \sigma(\pi - \theta)}$$

$A_{\text{fb}} = 0$ if charge symmetry conserved!

$np \rightarrow d\pi^0$ and $dd \rightarrow \alpha\pi^0$

Miller and van Kolck: For TRIUMF and IUCF χPT (**small p**)

$$A_{\text{fb}}(np \rightarrow d\pi^0) \propto \frac{\langle \pi^0 | H | \eta \rangle}{-0.0059 \text{ GeV}^2} - \frac{0.87}{\text{MeV}} \left(\delta M - \frac{\bar{\delta}M}{2} \right)$$

$$\sigma(dd \rightarrow \alpha\pi^0) \propto [A\langle \pi^0 | H | \eta \rangle + B\delta M + C\bar{\delta}M + D\langle \rho^0 | H | \omega \rangle]^2$$

where A, B, C, D are numbers and $m_n - m_p = \delta M + \bar{\delta}M$;

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$\delta M \sim \text{few MeV}$ — strong term

$-2 < \bar{\delta}M < -1 \text{ MeV}$ — EM term

$\delta M - \frac{1}{2}\bar{\delta}M \sim 2.1\text{--}3.6 \text{ MeV}$ for reasonable δM and $\bar{\delta}M$

$\langle \pi^0 | H | \eta \rangle = (-5900) \text{ to } (-4300) \text{ MeV}^2$, $g_\eta = ?$

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IUCF+TRIUMF constrain δM , $\bar{\delta}M$ and $\langle \pi^0 | H | \eta \rangle$ ($\leftrightarrow \beta_1$ in χPT)

Executive summary of Tuesday last week:

- Effectively one (1) data point (*s*-wave XS) PRL91, 142302 (2003).
- Formally leading amplitude suppressed (spin-isospin)
- **Short range amplitudes (higher orders) are significant**
- Many diagrams remain to be calculated
- Slow convergence
- dd wave function very complicated and largely unknown
- Realistic wfs overshoot data

AG *et al.*, PRC69, 044606 (2004)

Talks last week, available on-line soon!



$np \rightarrow d\pi^0$ should be easier!

- Both CS and CSB data exists!
- Wave functions well known
- Leading amplitudes not suppressed (cf. $pp \rightarrow pp\pi^0$)
- Plenty of work on CS $NN \rightarrow NN\pi^0$
- Influence of short-range needs to be checked

Theory for $NN \rightarrow NN\pi$

Extensive literature after IUCF (and TSL) $pp \rightarrow pp\pi^0$ expt.

Meyer et al., PRL 65, 2846 (1990); NPA 539, 633 (1992)

Bondar et al., PLB 356, 8 (1995)

Short-range correlations and HME

T-SH Lee and DO Riska, PRL 70, 2237 (1992).

CJ Horowitz, DK Griegel, and HO Meyer, PRC 49, 1337 (1994).

Off-shell amplitudes

E Hernandez and E Oset, PLB 350, 158 (1995). + many more

Modified χPT power counting scheme ($\sqrt{Mm_\pi}$ vs. m_π)

TD Cohen, JL Friar, GA Miller, and U van Kolck, PRC 53, 2661 (1996).

+ many, many more articles

Status summarized in C Hanhart, PRt 397, 155 (2004).

Cf. higher orders promoted for NN (Andreas on Tuesday)

CSB in $np \rightarrow d\pi^0$ II

Theory:

Cheung, Henley and Miller, PRL43, 1215 (1979)

Niskanen, Sebestyen, and Thomas, PRC38, 838 (1988)

$\eta-\pi^0$ mixing; $A_{fb} = -0.28\%$

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$\eta-\pi^0 + \chi\text{PT}$ (δM and $\bar{\delta}M$)

$A_{fb} = 0.23\% - 0.60\%$: Change of sign!

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Recent TRIUMF result confirms sign and size:

$A_{fb} = [0.172 \pm 0.080(\text{stat}) \pm 0.055(\text{syst})]\%$

Opper *et al.*, PRL 91 , 212302 (2003)

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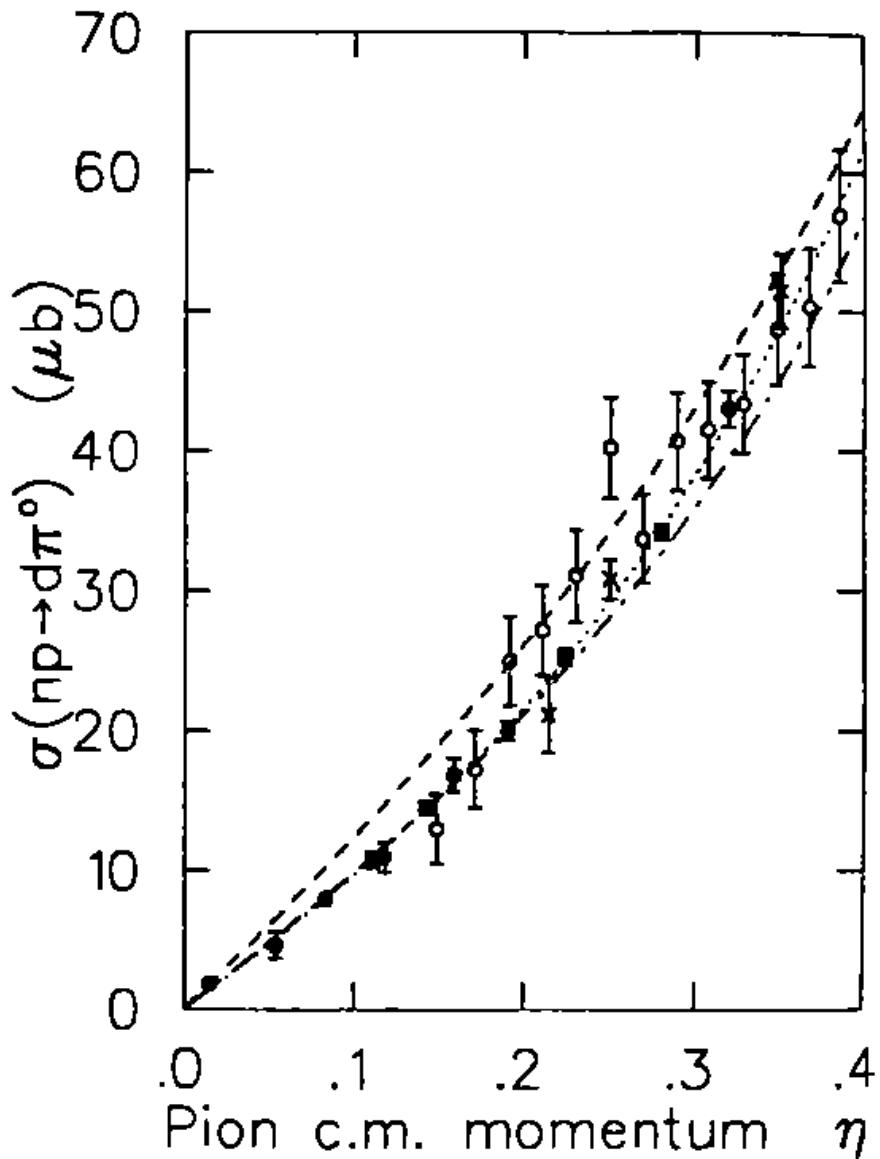
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No consistent chiral calculation carried out yet!

TRIUMF CS data for $np \rightarrow d\pi^0$ 

$$\eta = q/m_\pi$$

$\eta < 0.32$ for TRIUMF expt (•)

other data: $\frac{1}{2}\sigma(pp \rightarrow d\pi^+)_{CC}$

$$\sigma(np \rightarrow d\pi^0) = \frac{1}{2}(\alpha\eta + \beta\eta^3)$$

$$\alpha = 184 \pm 5 \mu b$$

$$\beta = 781 \pm 79 \mu b$$

Hutcheon et al., PRL64, 176 (1990); NPA535, 618 (1991).

Low-level analysis

CSB ($\alpha' \ll \alpha, \beta' \ll \beta$) \Rightarrow

$$\sigma(np \rightarrow d\pi^0) = \frac{1}{2}(\alpha\eta + \beta\eta^3) + \frac{1}{2}(\alpha'\eta + \beta'\eta^3)$$

Ignoring phases:

$$A_{\text{fb}}(np \rightarrow d\pi^0) = \frac{\sqrt{\alpha\beta'}\eta^2 + \sqrt{\alpha'\beta}\eta^2}{\alpha\eta + \beta\eta^3} \approx \sqrt{\frac{\beta'}{\alpha}}\eta + \frac{\sqrt{\alpha'\beta}}{\alpha}\eta$$

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A_{fb} depends on understanding of CS XS!

1st task: calculate $\sigma(np \rightarrow d\pi^0)$!



Partial wave analysis

Spin-isospin, parity and symmetry $\Rightarrow ({}^{2S+1}L_{np}l_\pi)$

$$^3P_1 s \left(\leftrightarrow \sqrt{\alpha}\right) \quad \text{CS} \quad ^1S_0 p, ^1D_2 p \left(\leftrightarrow \sqrt{\beta}\right)$$

$$^1P_1 s \left(\leftrightarrow \sqrt{\alpha'}\right) \quad \text{CSB} \quad ^3S_1 p, ^3D_1 p, ^3D_2 p \left(\leftrightarrow \sqrt{\beta'}\right)$$

$A_{\text{fb}} \leftrightarrow S\text{-}P\text{-wave singlet-singlet}$ and triplet-triplet interferences



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Need to calculate:

- CS triplet pion *s*-wave
 - CS singlet pion *p*-waves
 - CSB singlet pion *s*-wave
 - CSB triplet pion *p*-waves



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For now restrict to:

- ## • CS triplet pion *s*-wave

Power counting (CS)

Energy/momentum scales:

$$\mathbf{q} \lesssim \gamma = \sqrt{MB} \ll m_\pi < \Delta < P = \sqrt{Mm_\pi} \ll \Lambda_\chi$$
$$24 \lesssim 46 \ll 140 < 290 < 360 \ll 800 \text{ MeV}$$



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Pion production: new scale $P = \sqrt{Mm_\pi}$ (initial N momentum)

TD Cohen, JL Friar, GA Miller, U van Kolck, PRC 53, 2661 (1996).

Expand in $x \equiv \frac{P}{M} = \frac{m_\pi}{P} = \sqrt{\frac{m_\pi}{M}}$, not $\frac{m_\pi}{M}$

Also use η (mainly p -waves)

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With explicit Δ poorly separated scales. Looks really bad!

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Don't despair, I'll show later what saves the day...

CS Lagrangian

$$\begin{aligned}
 \mathcal{L}_{CS} = & -\frac{1}{4f_\pi^2} N^\dagger \boldsymbol{\tau} \cdot (\boldsymbol{\pi} \times \dot{\boldsymbol{\pi}}) N \\
 & + \frac{g_A}{2f_\pi} \left\{ N^\dagger \textcolor{green}{\boldsymbol{\tau}} \cdot \textcolor{red}{\vec{\sigma}} \cdot \vec{\nabla} \boldsymbol{\pi} N - \frac{1}{2m_N} \left[iN^\dagger \textcolor{green}{\boldsymbol{\tau}} \cdot \dot{\boldsymbol{\pi}} \textcolor{red}{\vec{\sigma}} \cdot \vec{\nabla} N + \text{H.c.} \right] \right\} \\
 & + \frac{h_A}{2f_\pi} \left\{ N^\dagger \boldsymbol{T} \cdot \vec{S} \cdot \Delta(\vec{\nabla} \boldsymbol{\pi}) + \text{H.c.} \right. \\
 & \left. - \frac{1}{M} \left[iN^\dagger \boldsymbol{T} \cdot \dot{\boldsymbol{\pi}} \vec{S} \cdot \vec{\nabla} \Delta + \text{H.c.} \right] \right\} + \dots
 \end{aligned}$$

$\Rightarrow \pi$ prod operator \mathcal{O}



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$\Rightarrow \pi$ prod operator \mathcal{O}

Sandwiched between wfs (we'll use a menagerie of wfs):

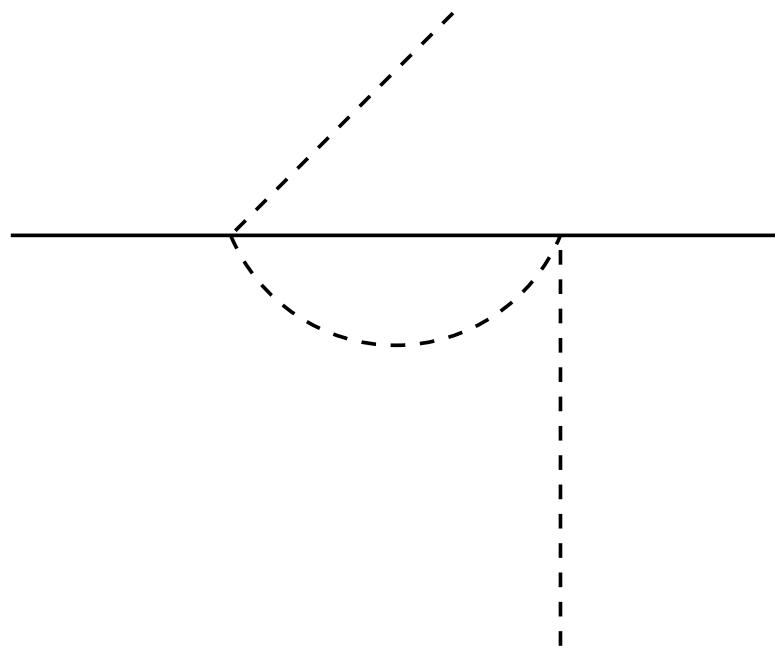
$$\mathcal{M} \sim \langle f | \mathcal{O} | i \rangle$$



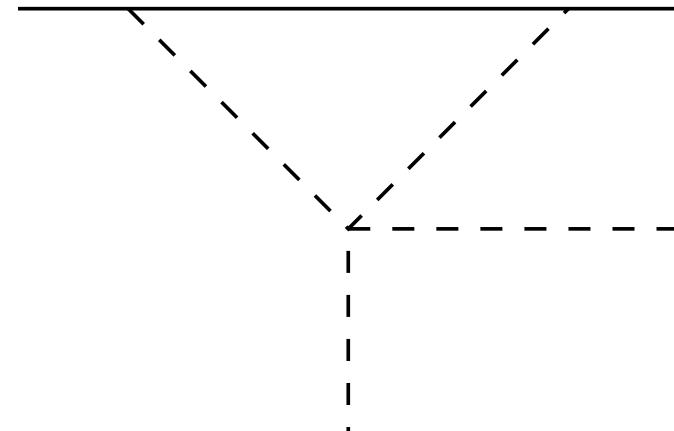
Power counting in loops

Two cases:

- P can be carried by nucleon: scales with m_π (left)
- P has to pass through a pion: scales with P (right)

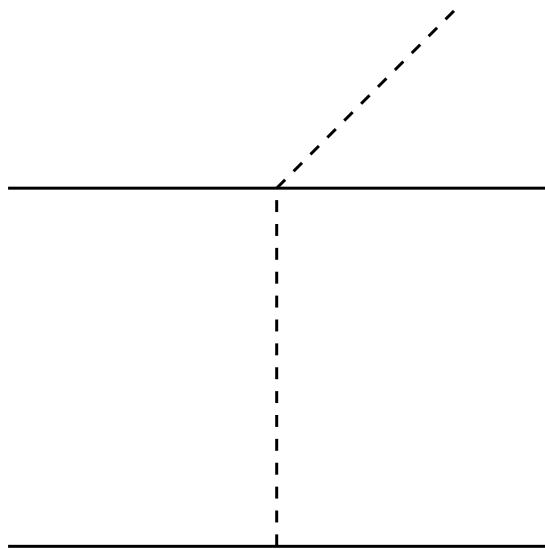


P-reducible

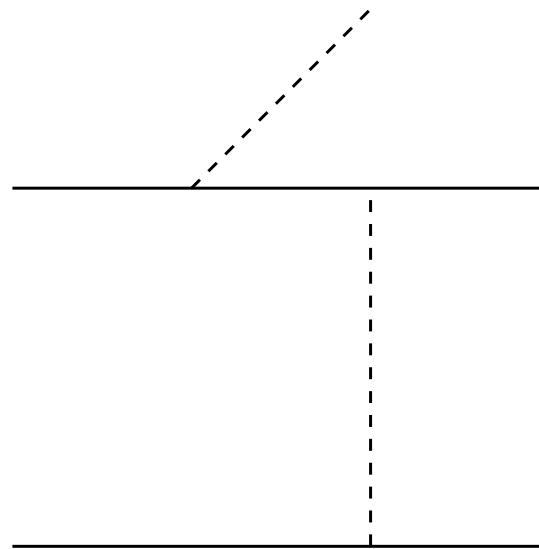


P-irreducible

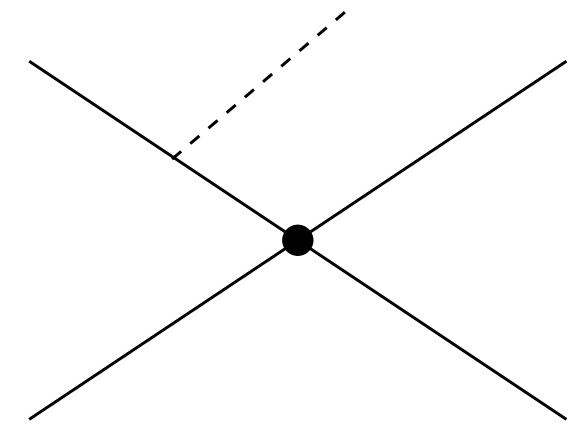
CS s -wave LO, $O(x)$



(a)



(b)

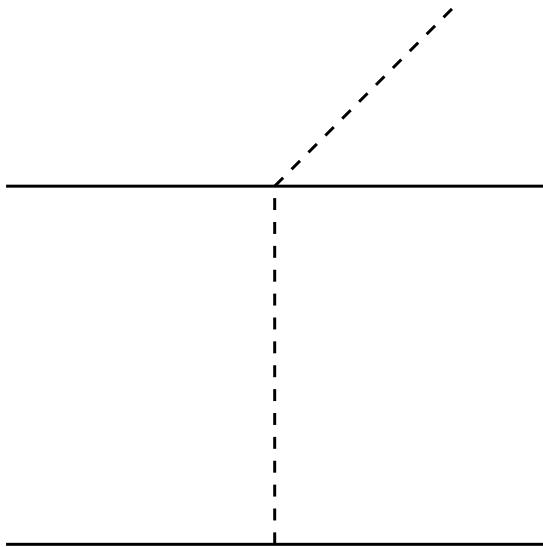


(c)

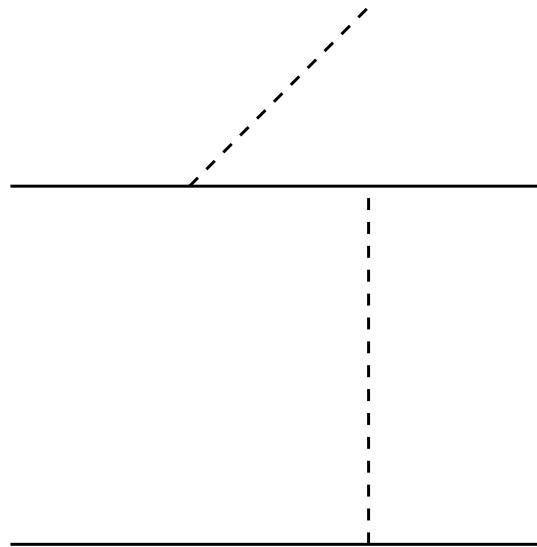
Weinberg-Tomozawa (a) + single-nucleon π production (b,c)

Pion exchange and contact term \rightarrow wave function

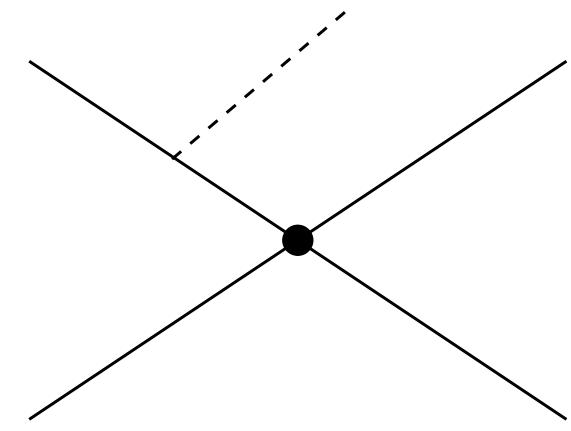
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(b)



(c)

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Caveats for W-T and 1-body

Remember the separation:

$$\sigma = \frac{1}{2}(\alpha\eta + \beta\eta^3)$$



CAUTION!

Remember the separation:

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Weinberg-Tomozawa $\propto \omega \approx \mu + \frac{q^2}{2\mu} \Rightarrow$

CAUTION!

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$$\sigma = \frac{1}{2}[\alpha_0\eta + (\alpha' + \beta)\eta^3]$$

Leakage from α to β !

Experiments with ang dis or polarization needed!

What is a one-body term, really?

Power counted using OPE, leads to ‘off-shell’ nucleon

But: wfs (potential or chiral) have ‘on-shell’ nucleons and fixed kinematics (FK) pion exchanges

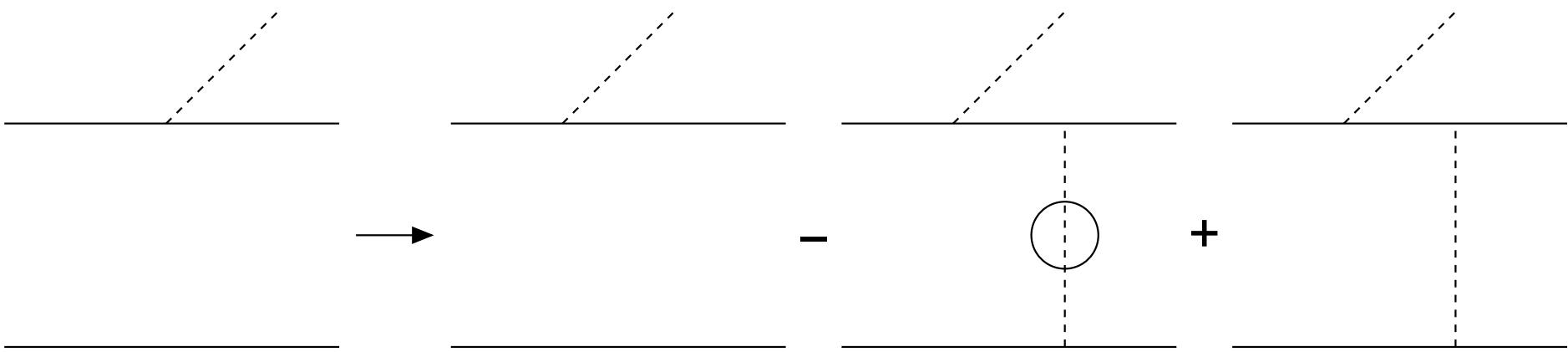
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Correction: 1B - FKOPE + full OPE (TOPT), in ISI and FSI

For FSI (first correction):

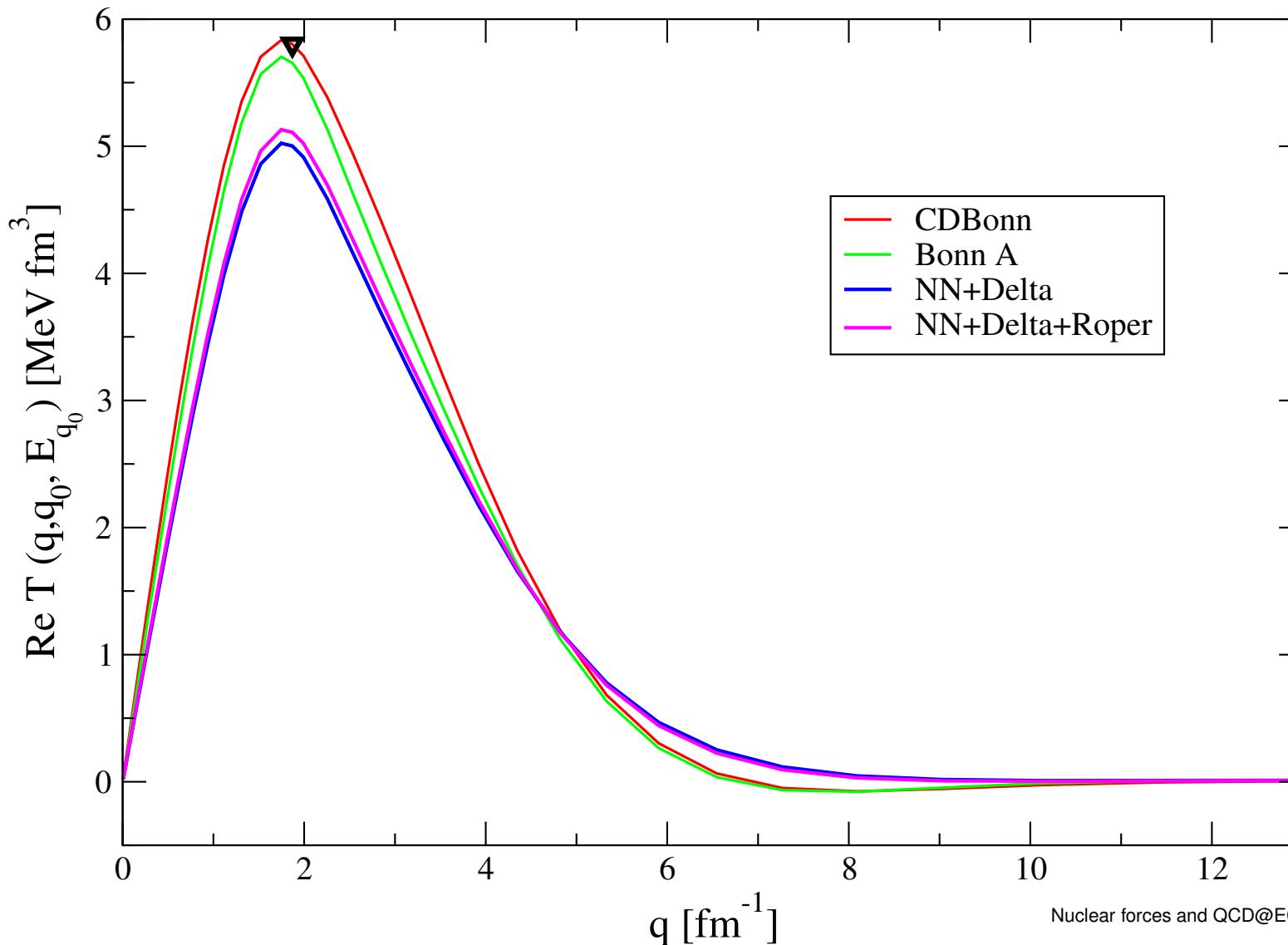


Numerical investigation needed!

Hopefully small effect \Rightarrow further corrs not needed?

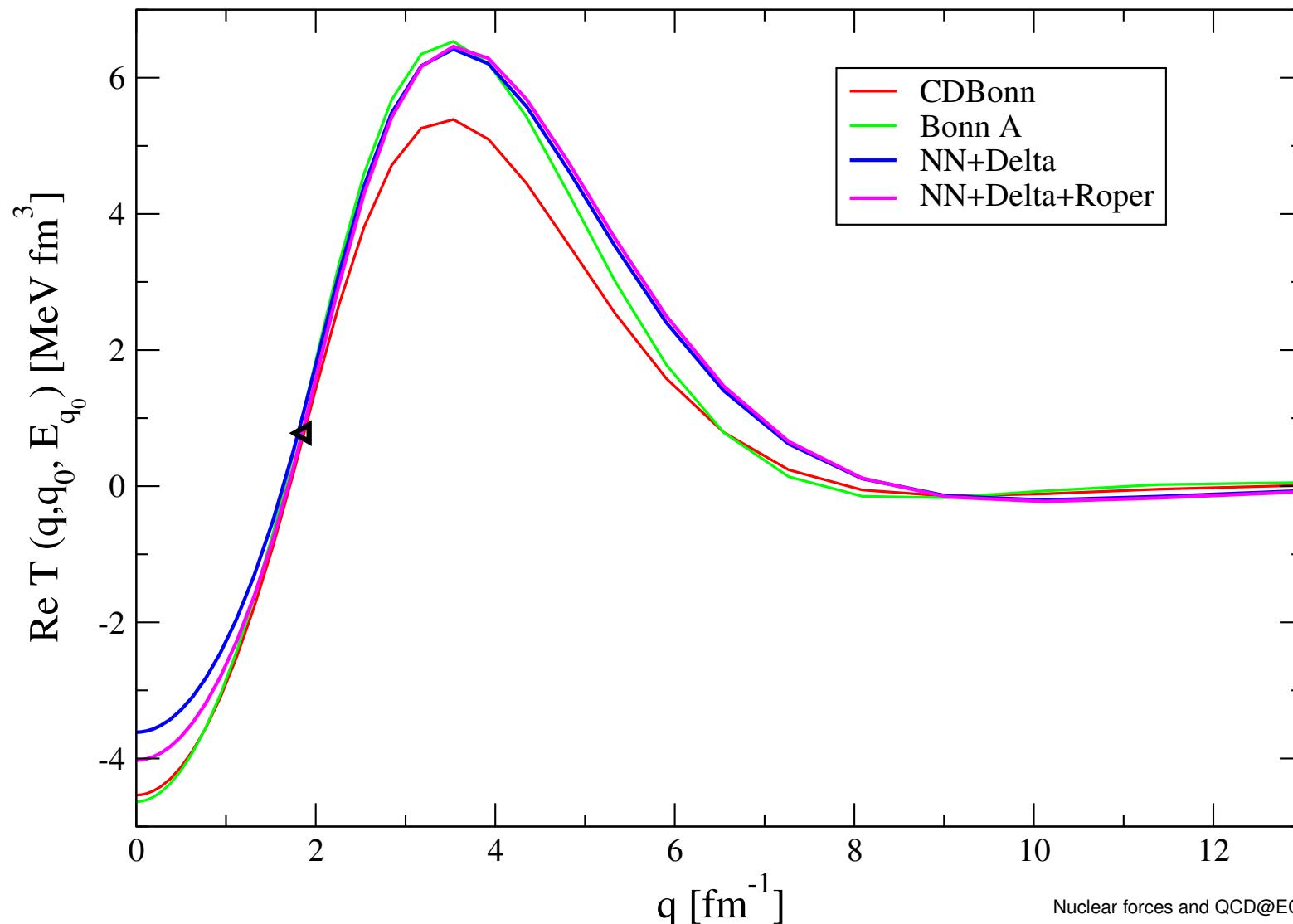
Half-off-shell $T(^3P_1)$

3P_1 half-shell T-matrix ($E_{\text{lab}}=290 \text{ MeV}$)

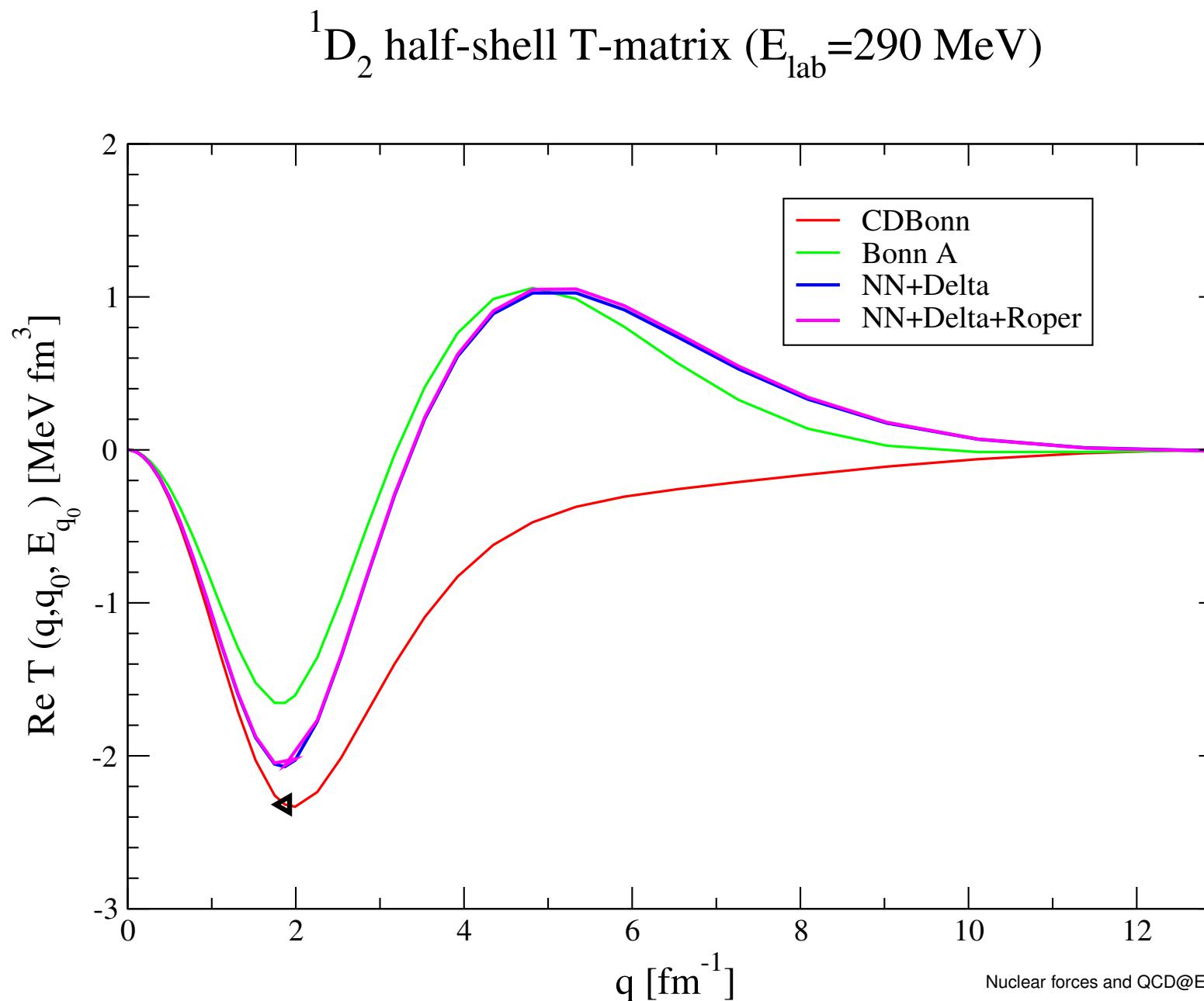


Half-off-shell $T(^1S_0)$

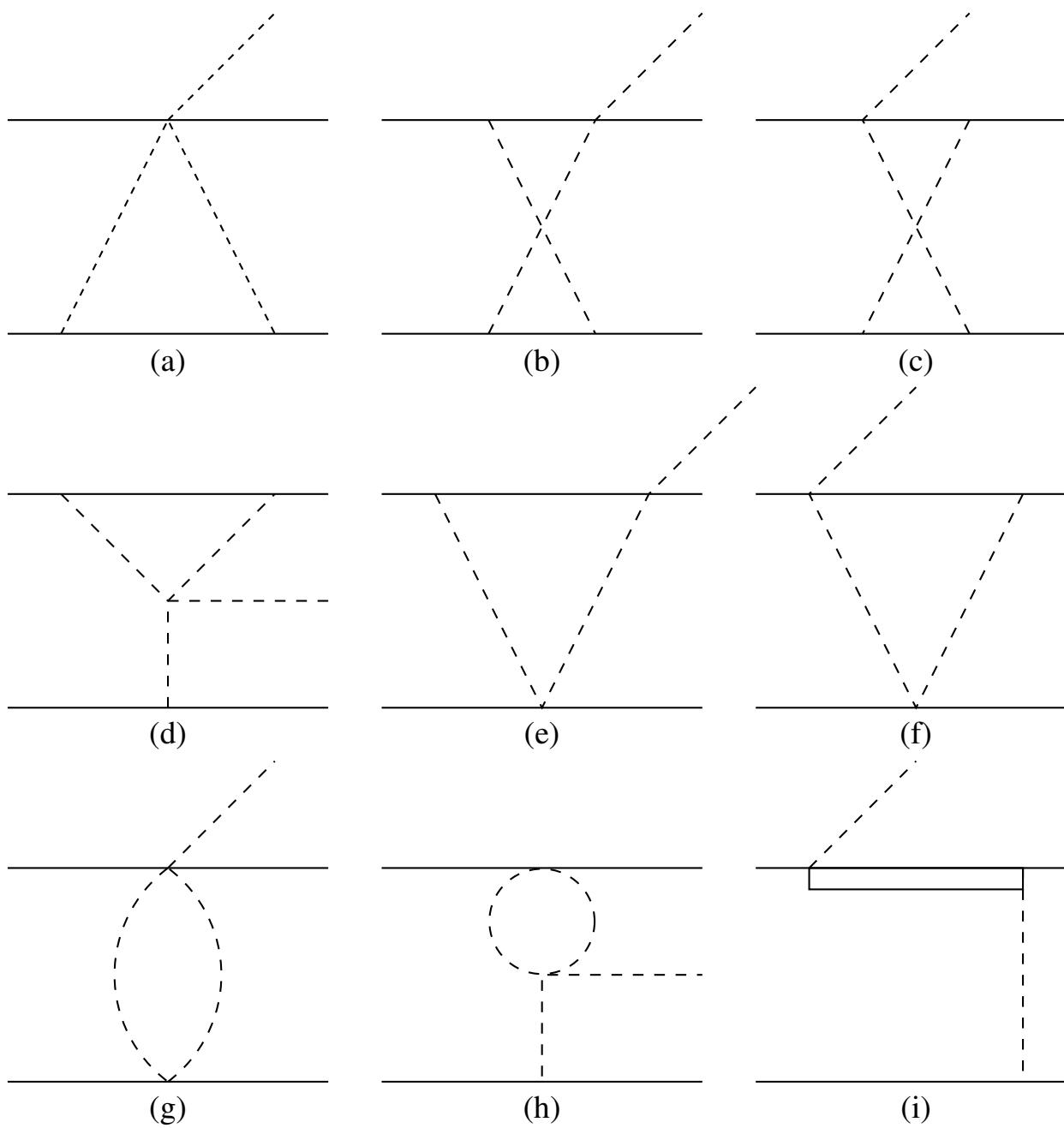
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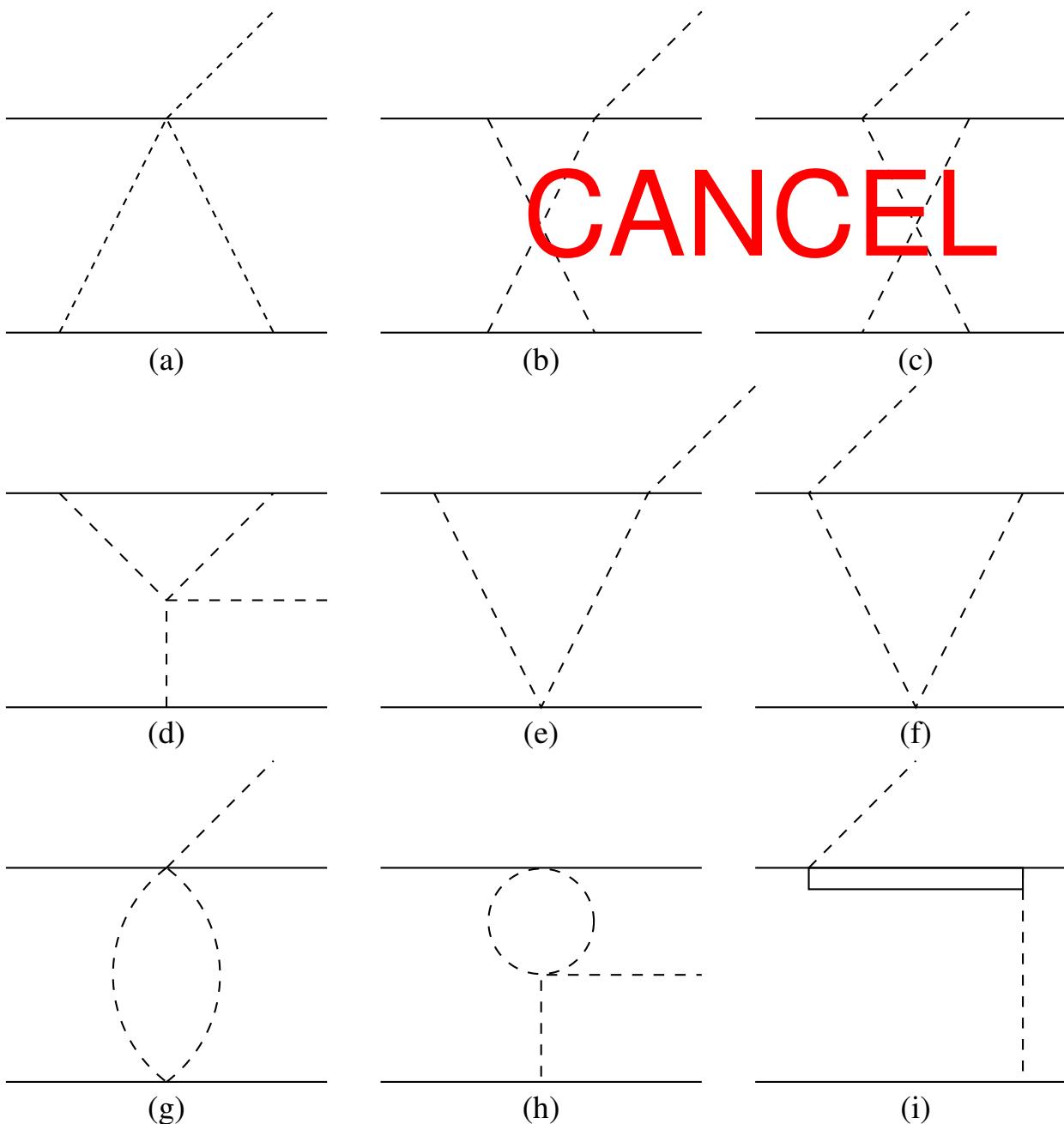
Half-off-shell $T(^1D_2)$



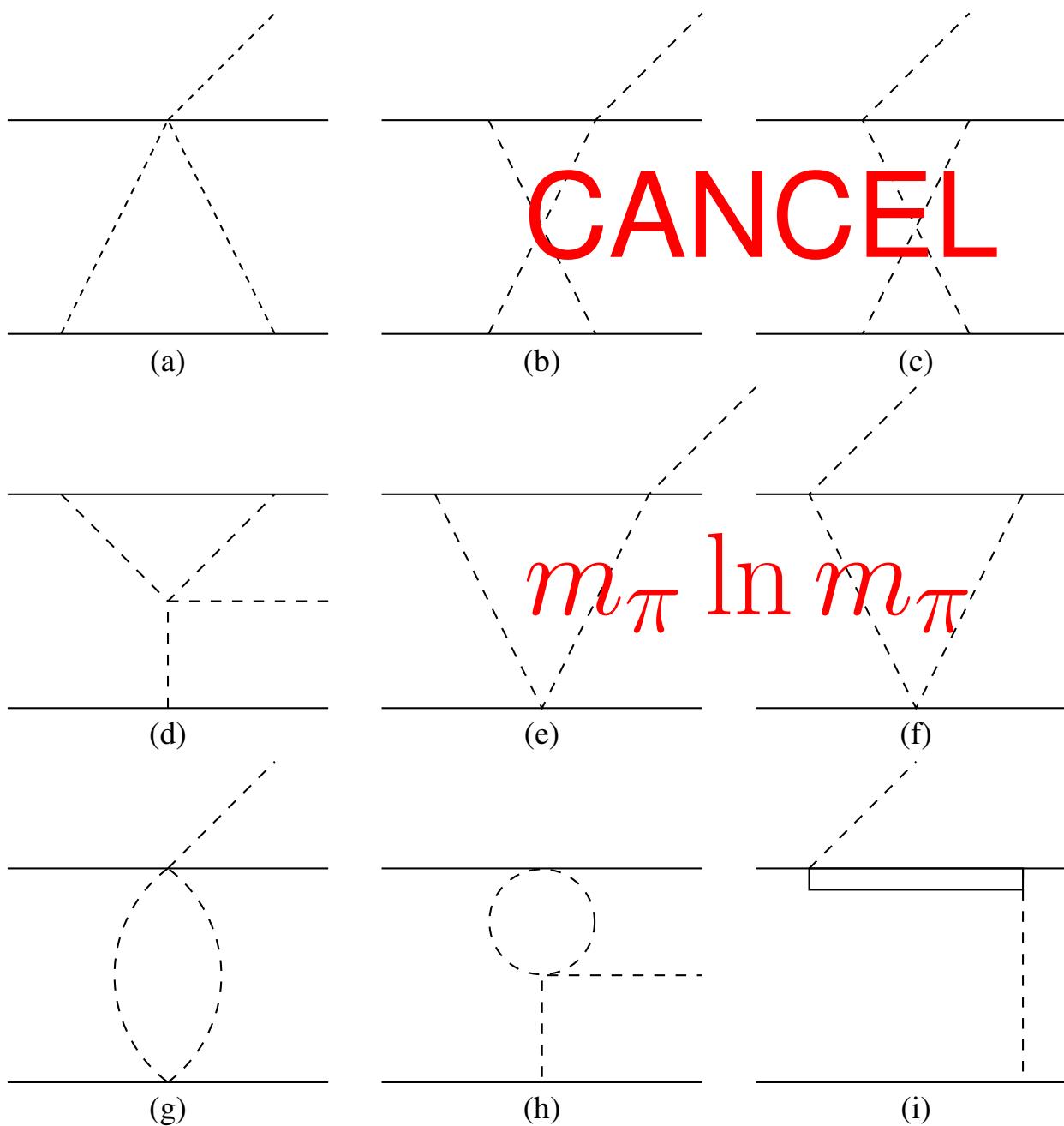
Hanhart & Kaiser, [PRC66, 054005 (2002)]



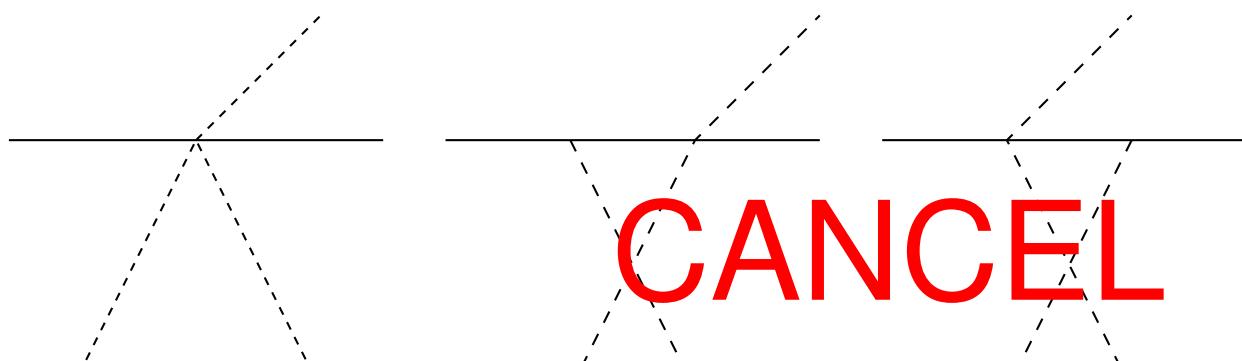
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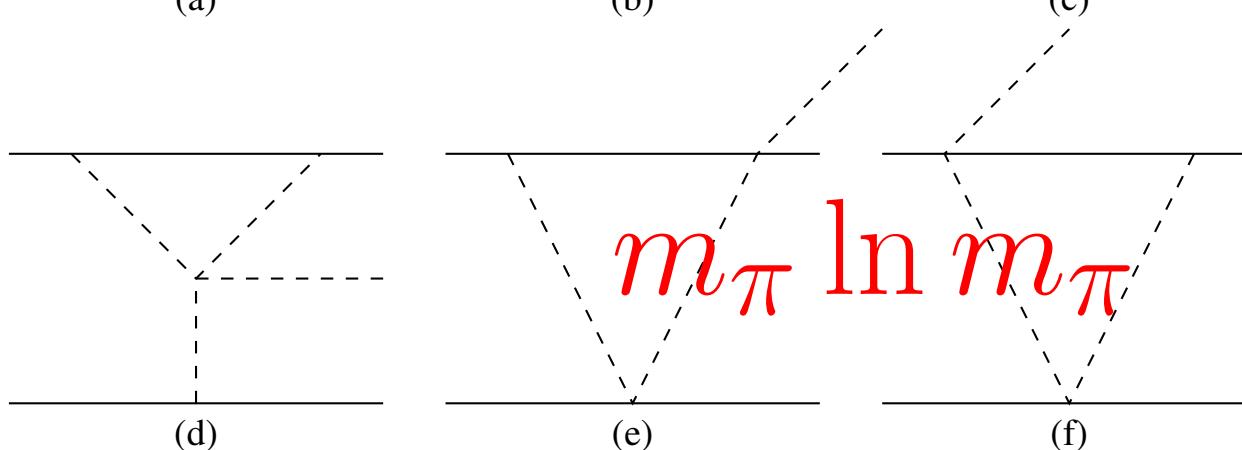
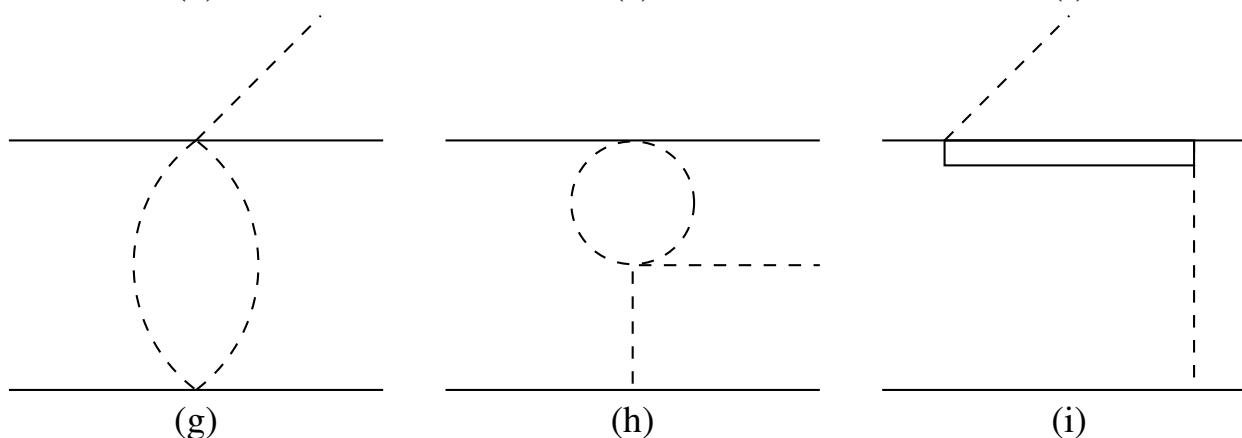
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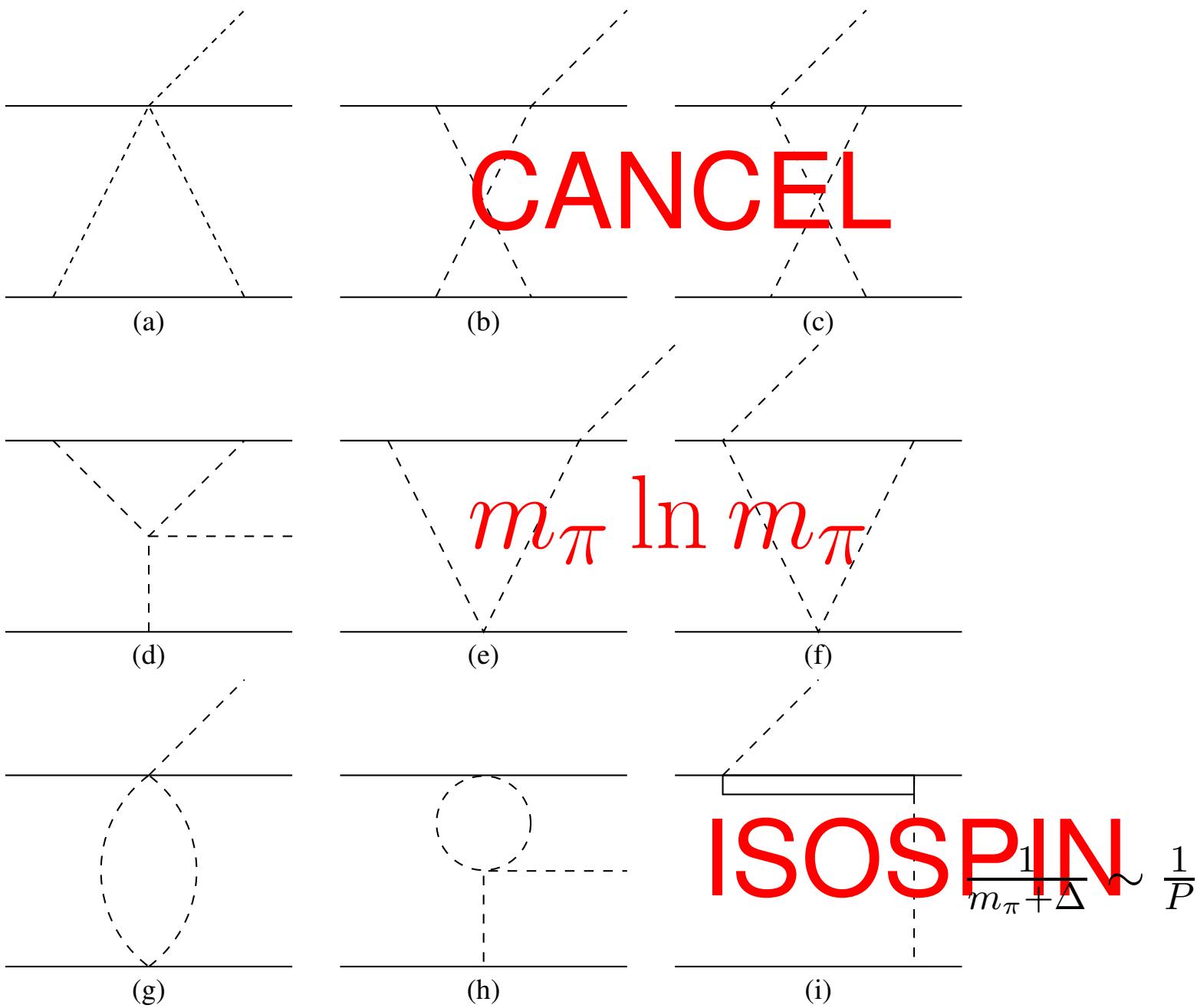


CANCEL

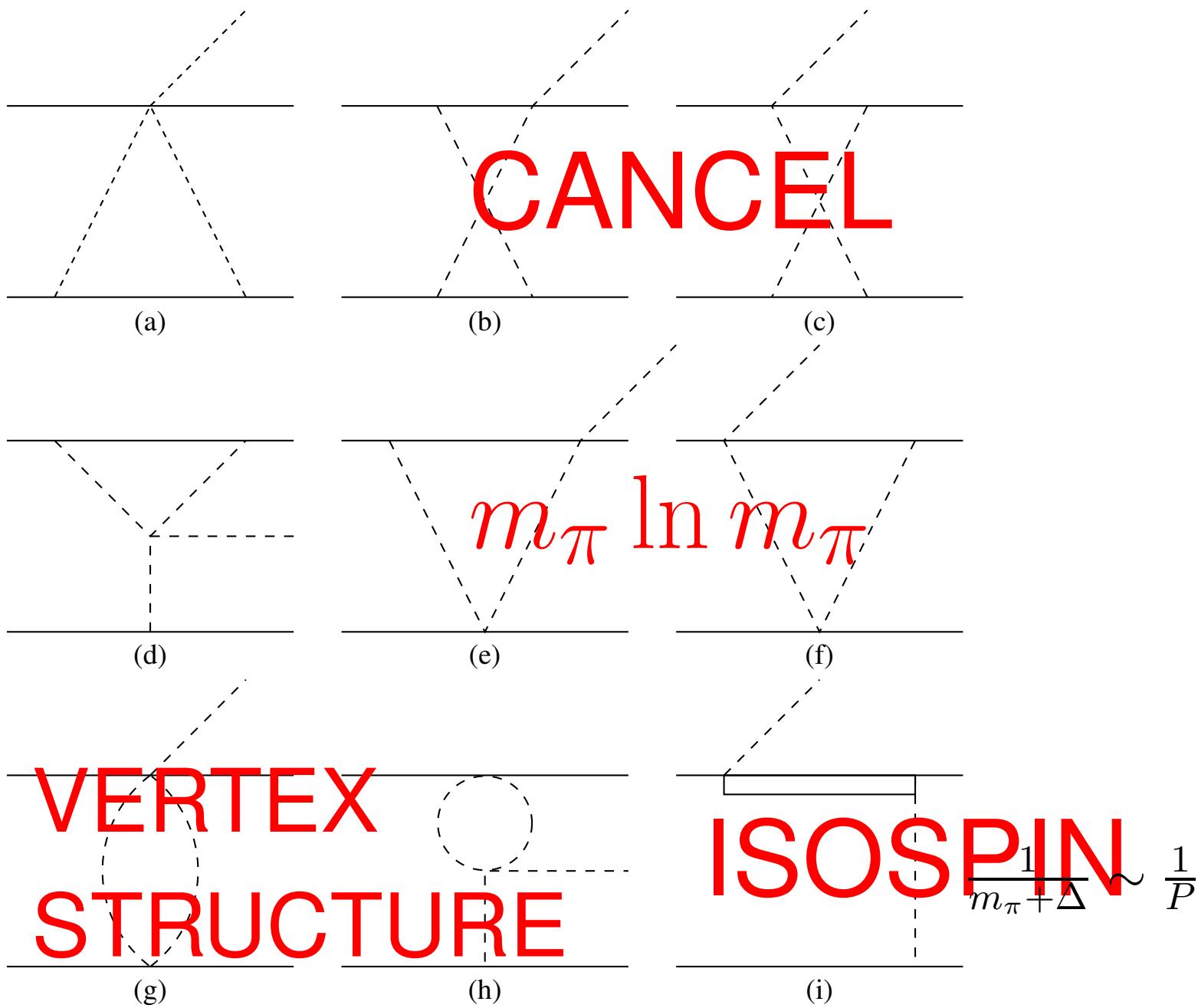
 $m_\pi \ln m_\pi$ 

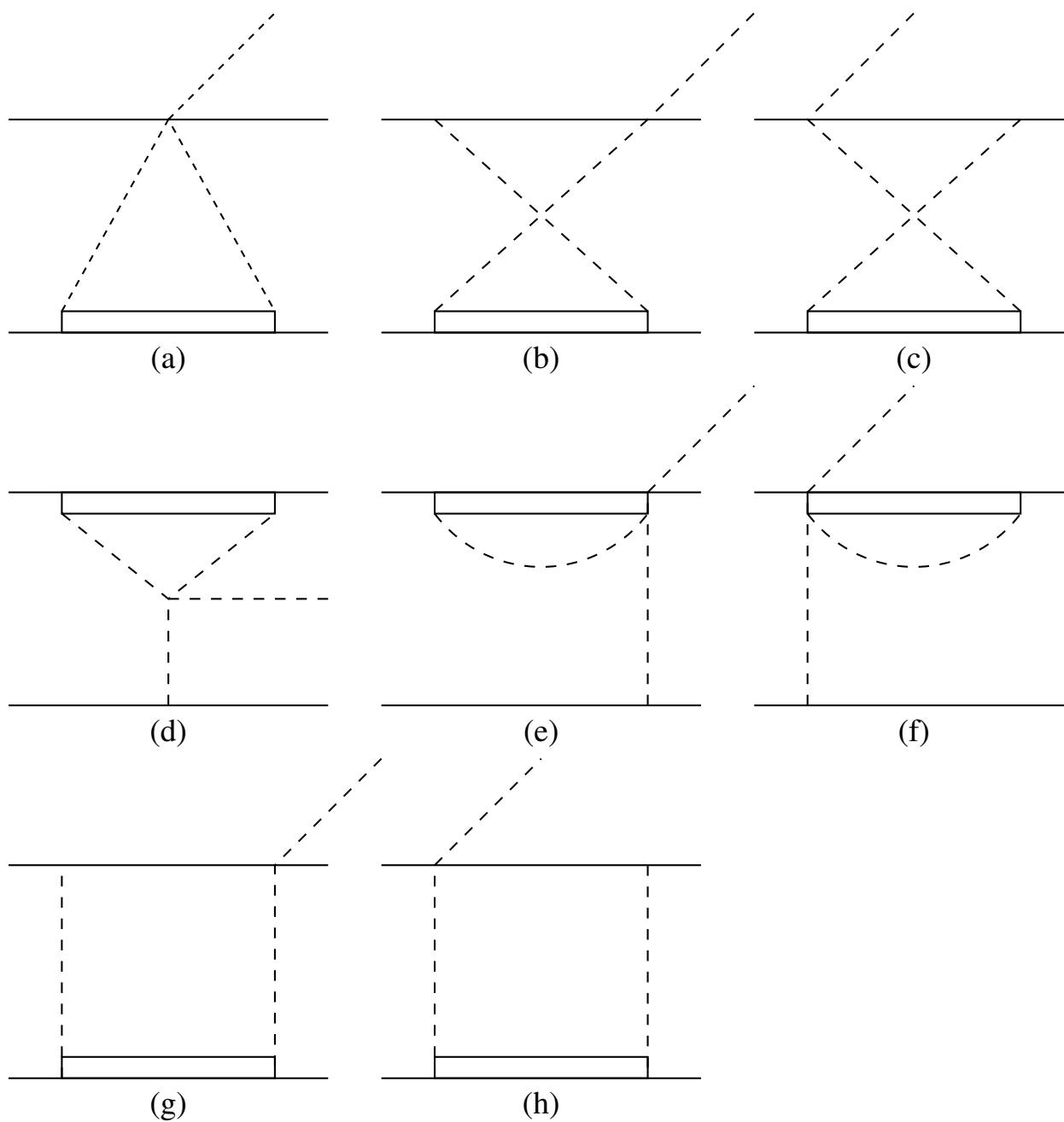
$$\frac{1}{m_\pi + \Delta} \sim \frac{1}{P}$$

Hanhart & Kaiser, [PRC66, 054005 (2002)]

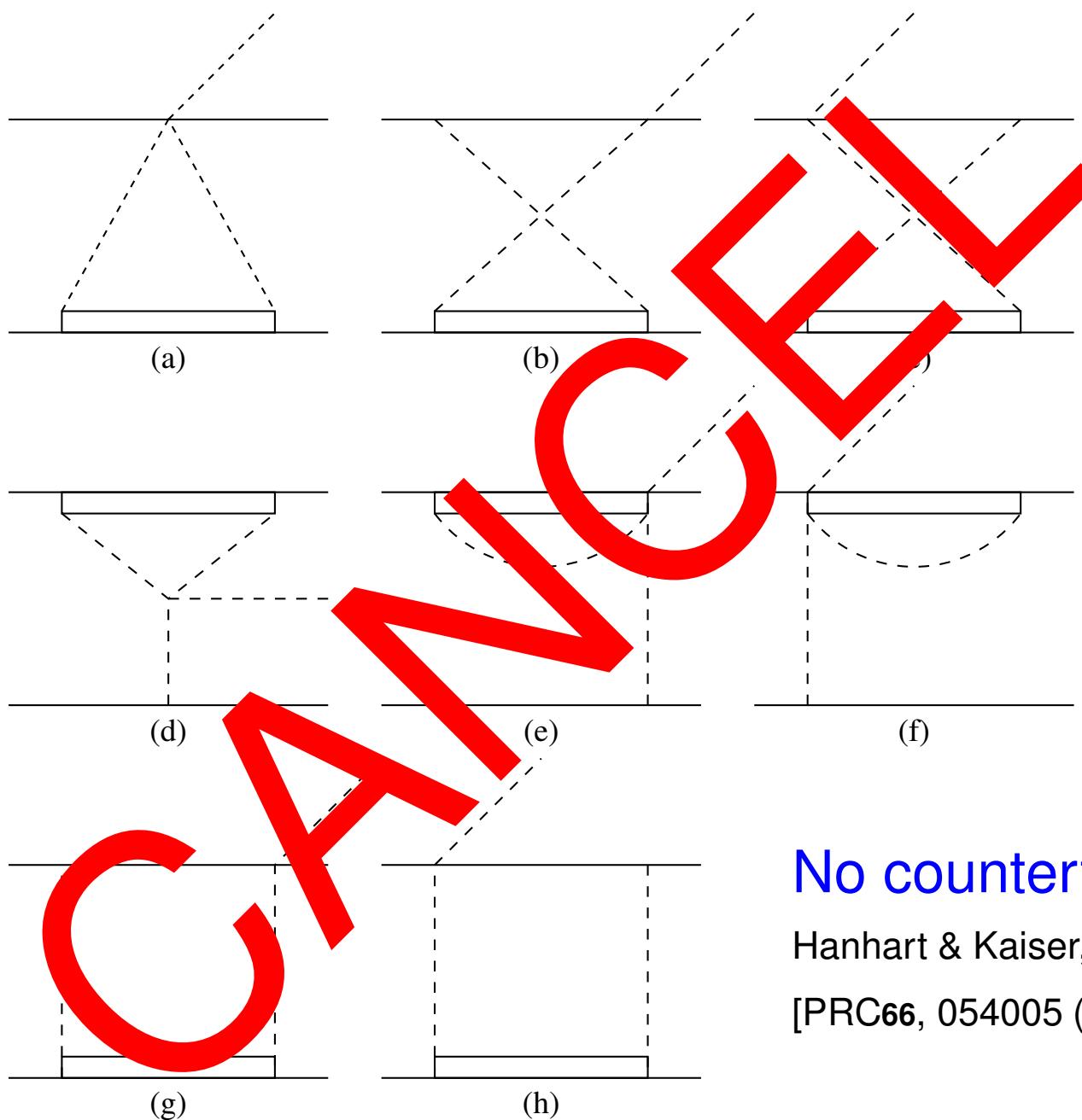


Hanhart & Kaiser, [PRC66, 054005 (2002)]



CS *s*-wave cont'd

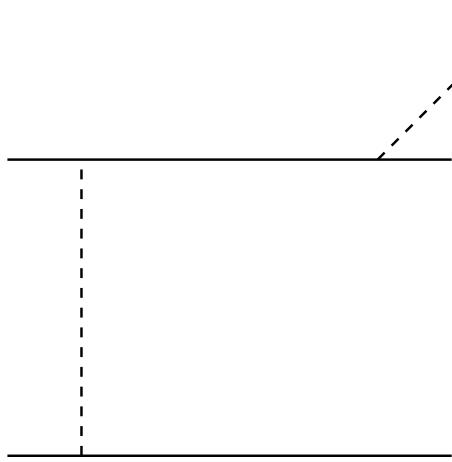
CS *s*-wave cont'd



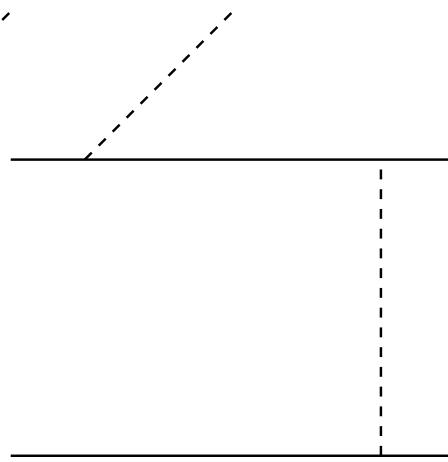
No counterterm

Hanhart & Kaiser,

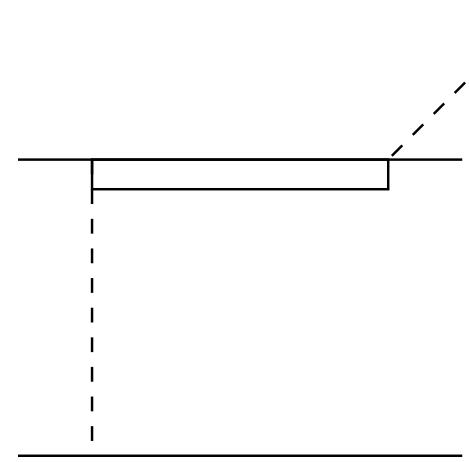
[PRC66, 054005 (2002)]

CS p -wave LO, $O(\eta)$ 

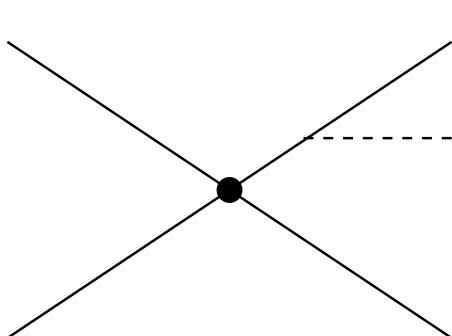
(a)



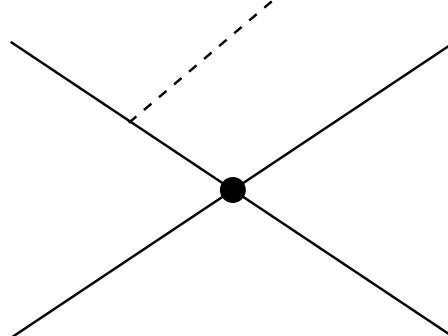
(b)



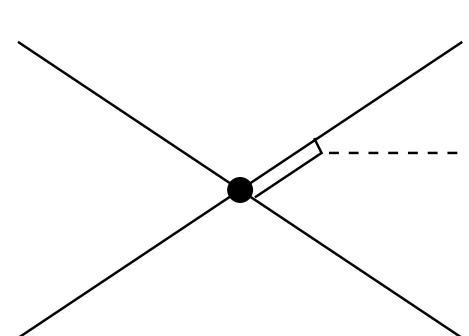
(c)



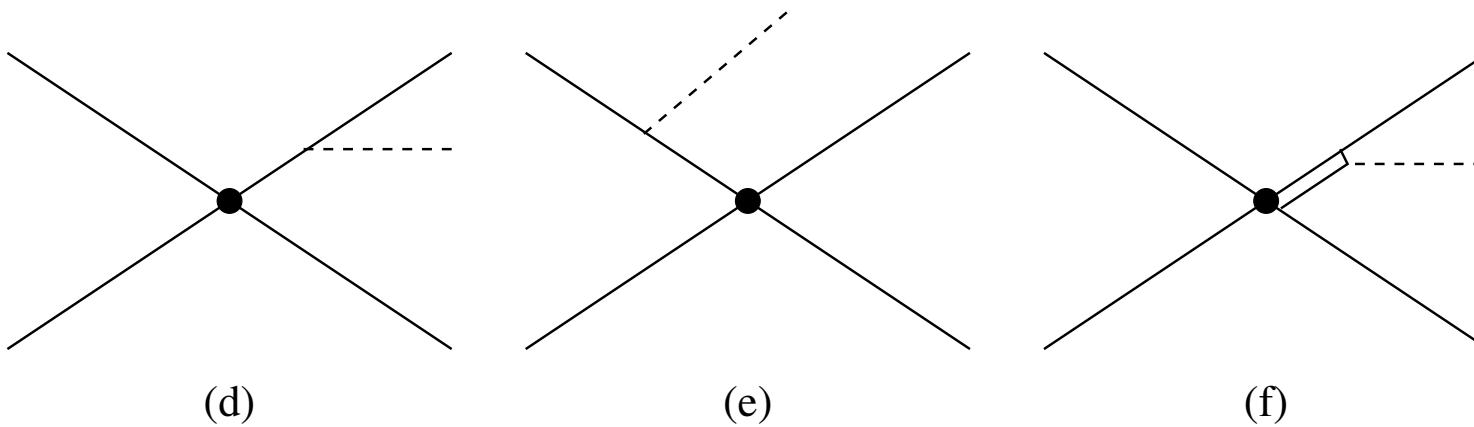
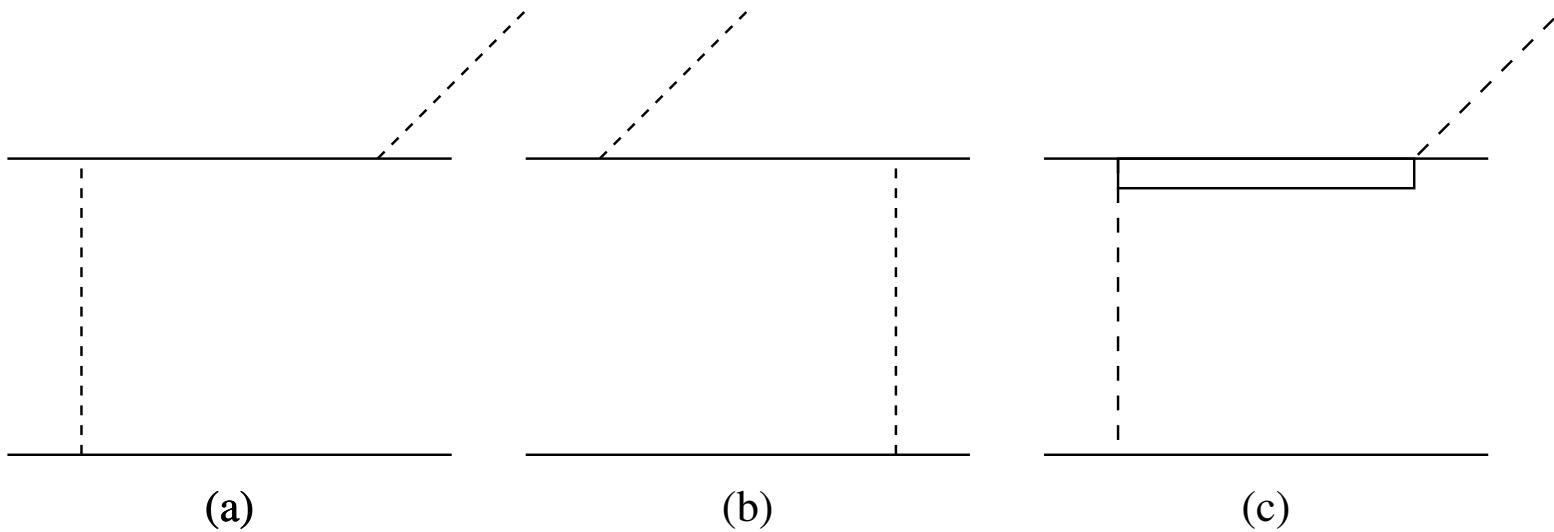
(d)



(e)



(f)

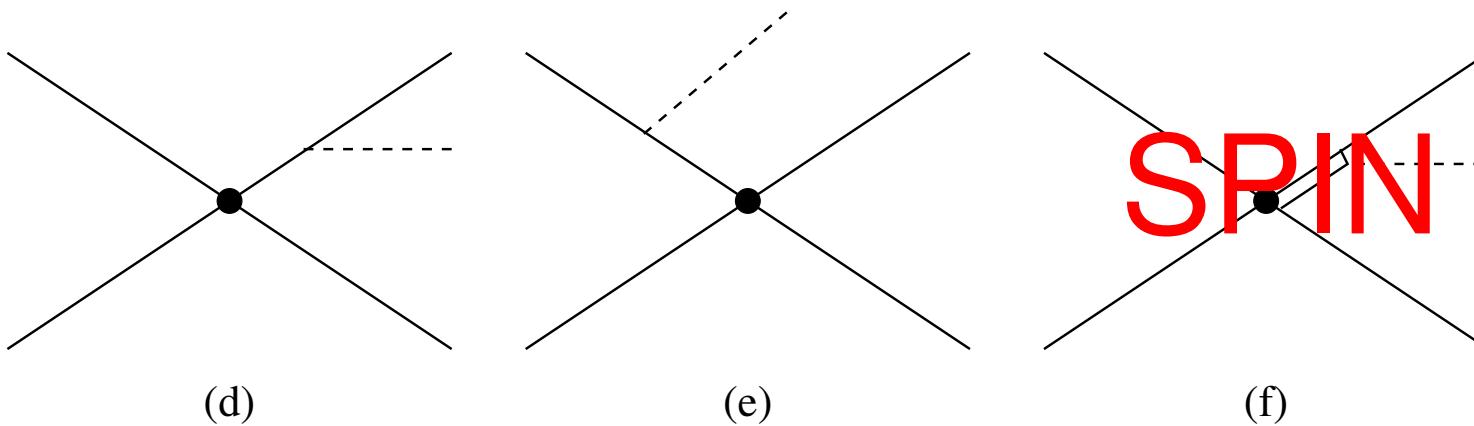
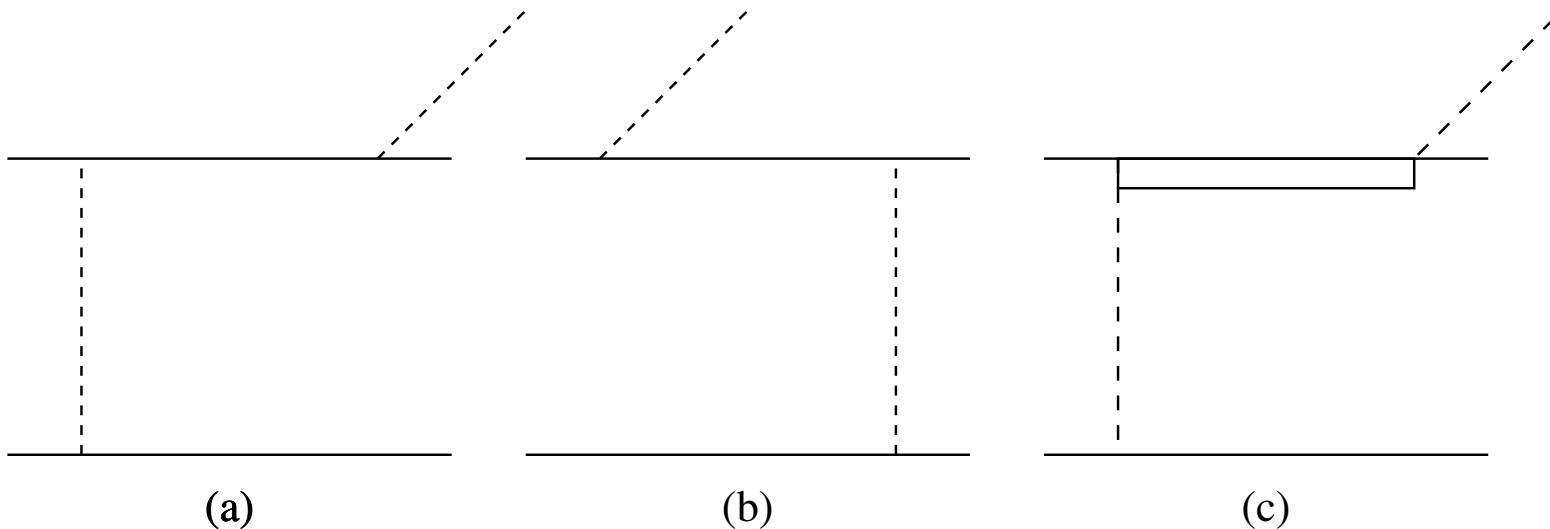
CS p -wave LO, $O(\eta)$ 

Δ propagator: $1/(\Delta - m_\pi) \sim 1/m_\pi$

No surviving Δ scale!



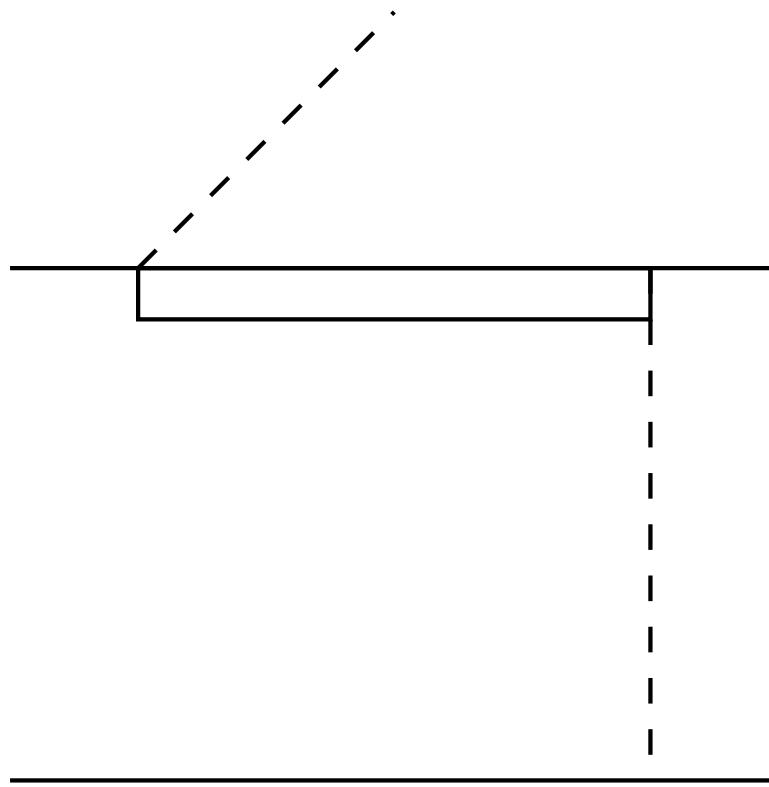
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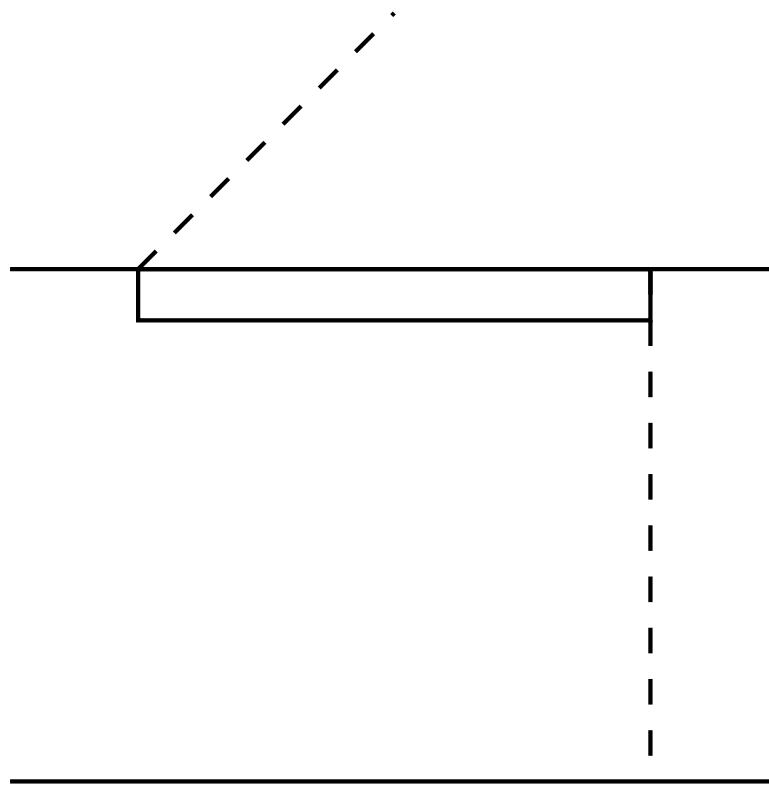
No surviving Δ scale!

CS p -wave NLO, $O(x\eta)$



Δ propagator: $1/(\Delta + m_\pi) \sim 1/P$

Δ scale disappears again!

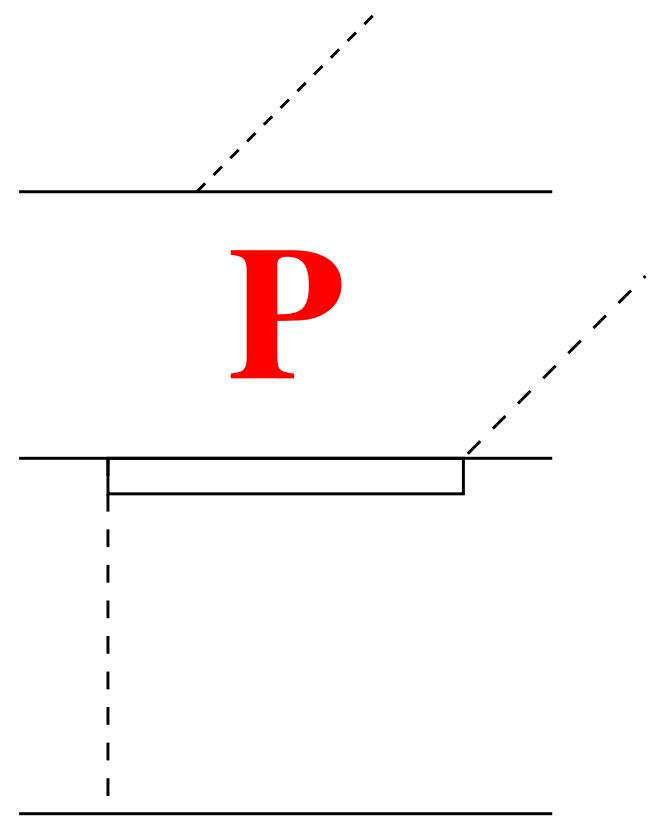
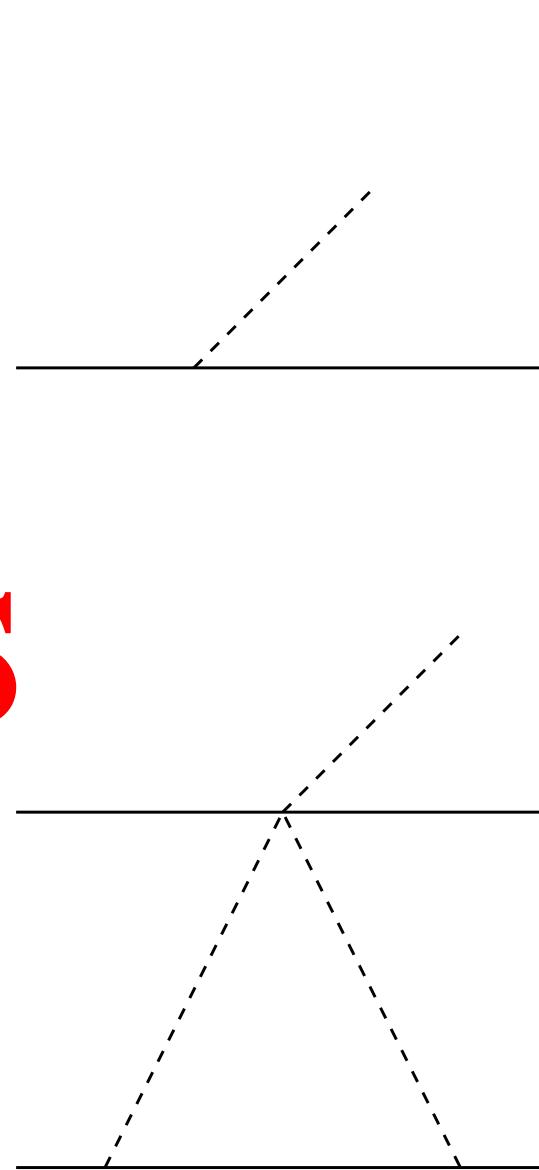
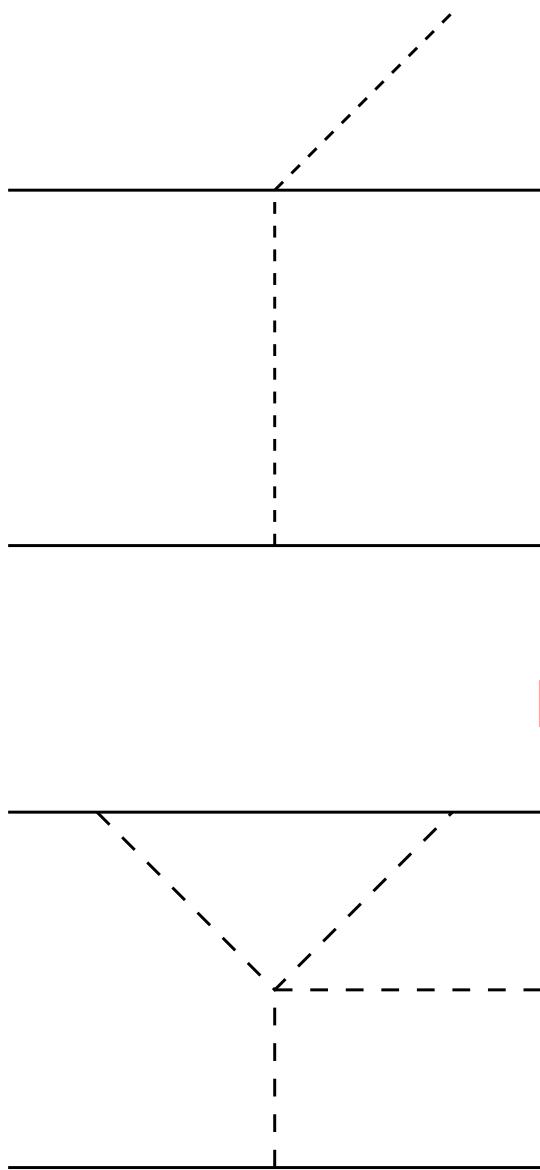


Δ propagator: $1/(\Delta + m_\pi) \sim 1/P$

Δ scale disappears again!

Deuteron isoscalar \Rightarrow no NLO p -wave diagram!

Summary CS diagrams to NLO



Calculate α to NLO

Impl. of WT, one-body and Δ (p -wave) straight-forward

Hanhart & Kaiser, PRC66, 054005 (2002) surviving DR'd loops:

$$(\text{isospin}) \times \frac{g_A^3 p(\boldsymbol{\sigma}_1 + \boldsymbol{\sigma}_2) \cdot \mathbf{p}}{256 f_\pi^5}$$

In r -space, with gaussian cut-off Λ :

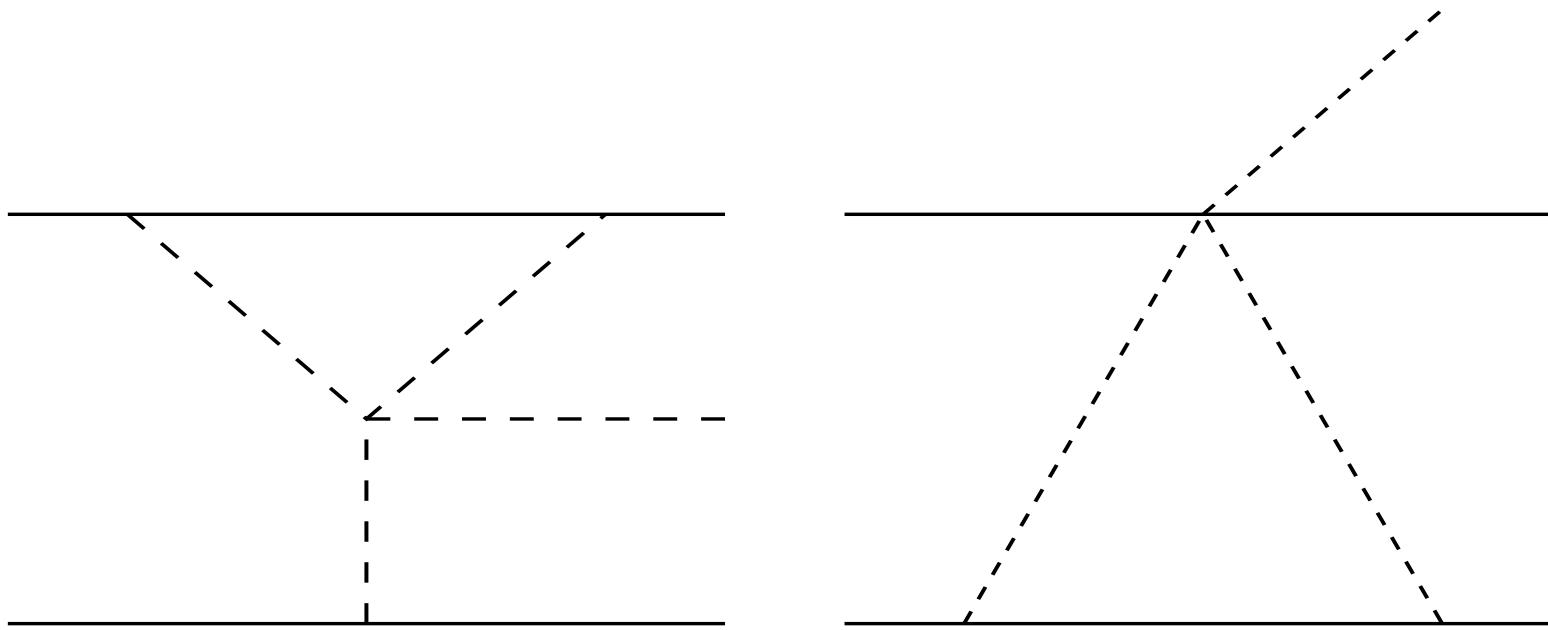
$$\mathcal{M} \propto \Lambda^5 \int dr u(r) F(\Lambda r) v(r)$$

$$F(x) = \frac{1}{x} - x + \sqrt{2} \left(\frac{1}{x^2} + 2 - x^2 \right) e^{-\frac{x^2}{2}} \int_0^{\frac{x}{\sqrt{2}}} dt e^{t^2} \rightarrow \frac{8}{x^5} \text{ (for large } x\text{)}$$

$$\Rightarrow \mathcal{M} \propto \Lambda \text{ (for large } \Lambda\text{)}$$

Renormalization I (operators)

Renormalize div. loop diagrams (using DR)

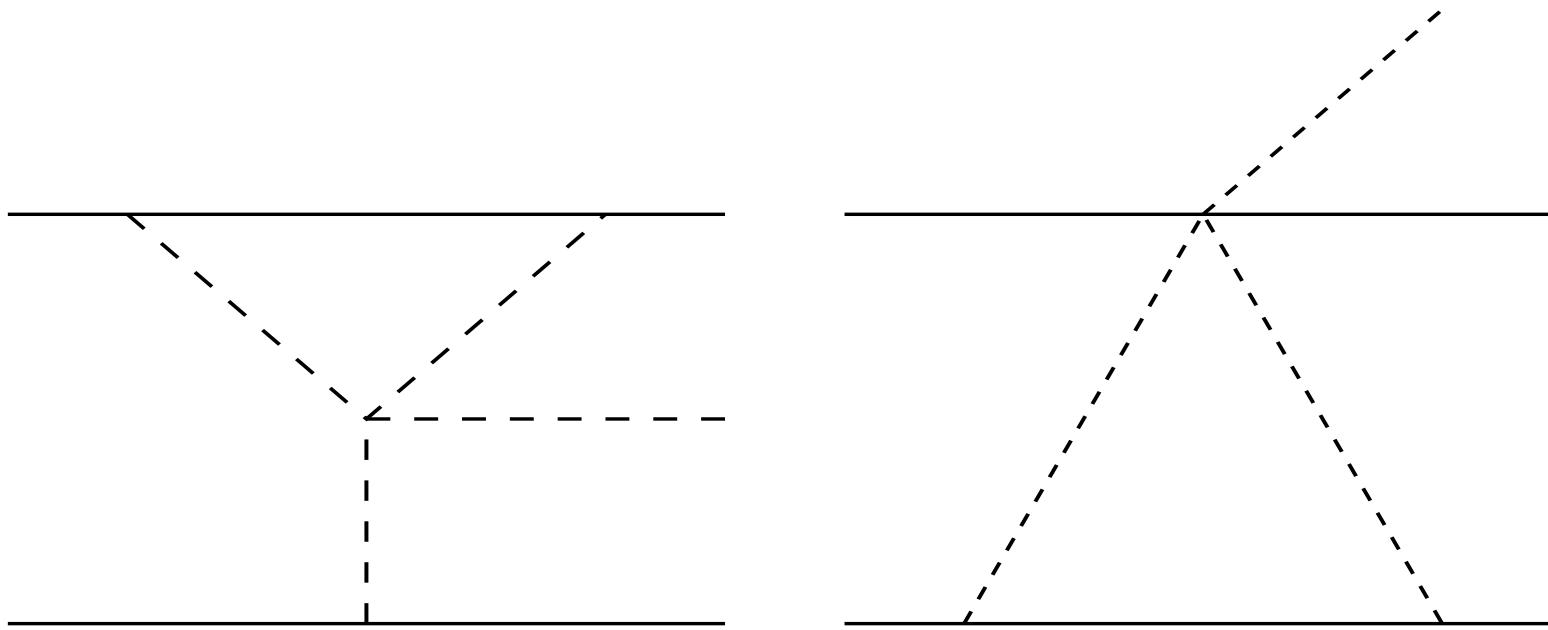


DR \Rightarrow divergencies removed from operators, excellent!

Renormalization I (operators)



Renormalize div. loop diagrams (using DR)



DR \Rightarrow divergencies removed from operators, excellent!

But we need MEs and XSs, and thus wfs

Might not be convergent if operator is singular

Breaks chiral sym! (linear div)



First results (preliminary)

Without Δ amplitudes:

| Λ [MeV] | α [μb] | | | β [μb] |
|-----------------|----------------------------|-----|------|---------------------------|
| | 500 | 800 | 1200 | |
| NijmI | 333 | 293 | 259 | 723 |
| exp't | $184 \pm 5 \mu\text{b}$ | | | $781 \pm 79 \mu\text{b}$ |

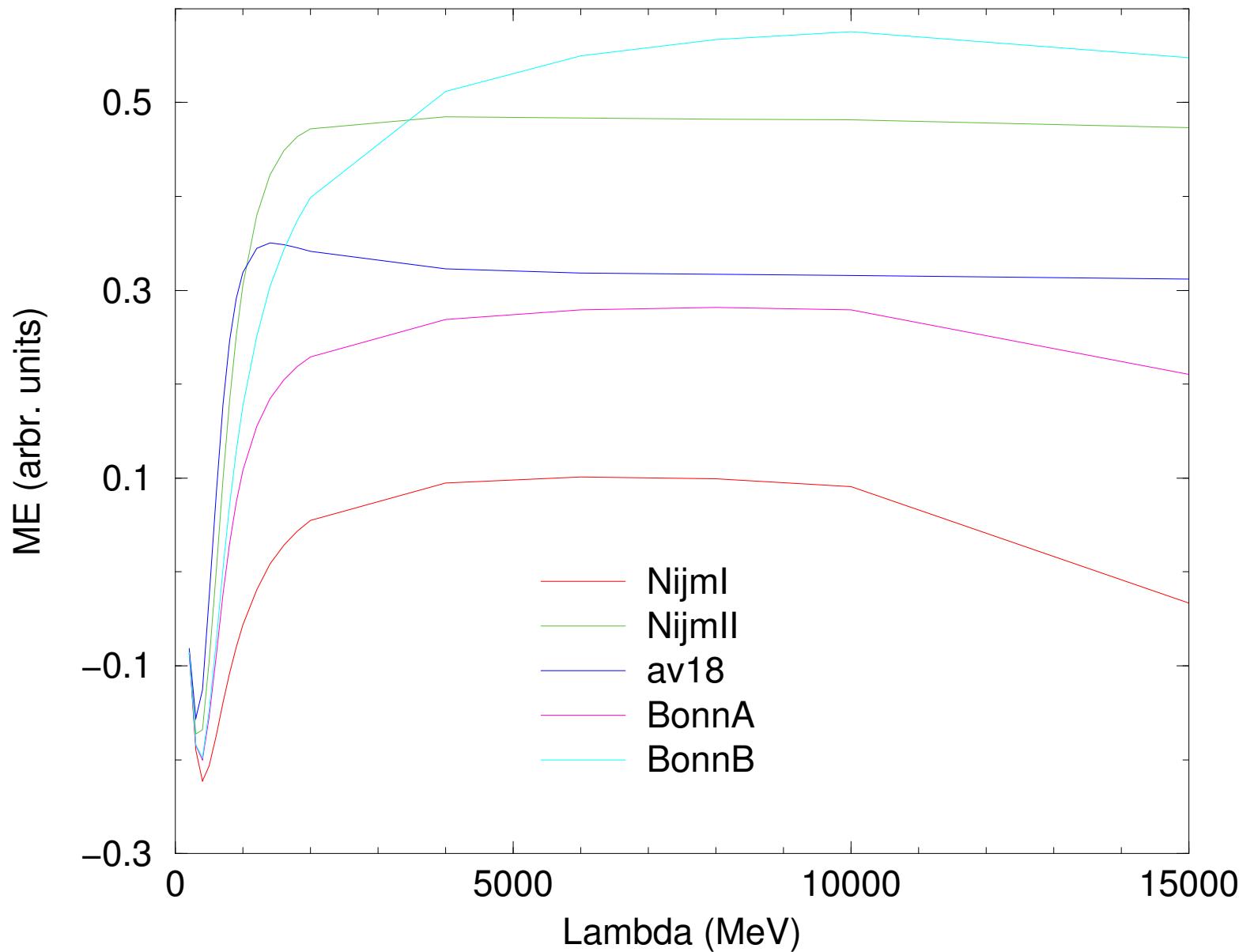
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|-----------------|----------------------------|------|--------------------------|---------------------------|
| | 500 | 800 | 1200 | |
| NijmI | 333 | 293 | 259 | 723 |
| NijmII | 217 | 134 | 86.9 | 653 |
| av18 | 146 | 80.0 | 61.2 | 662 |
| BonnA | 253 | 190 | 153 | 604 |
| BonnB | 265 | 190 | 137 | 653 |
| exp't | $184 \pm 5 \mu\text{b}$ | | $781 \pm 79 \mu\text{b}$ | |

What's going on here?

NLO loops



Renormalization II (\mathcal{M})

Renormalize ME when ISI and FSI (non-pert wfs) included

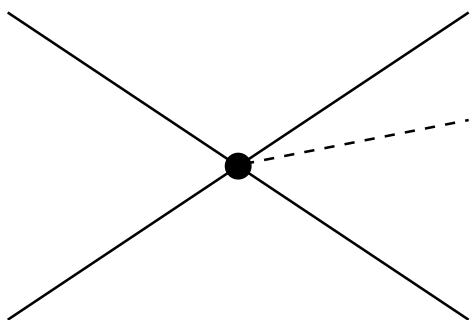
Renormalization II (\mathcal{M})



Renormalize ME when ISI and FSI (non-pert wfs) included

Regularize non-pert wfs???

2nd ren num, counter term to restore (?) χ symmetry:



$$(\text{isospin}) \times \frac{\tilde{d}_0(\Lambda)}{M f_\pi^3} (\boldsymbol{\sigma}_1 + \boldsymbol{\sigma}_2) \cdot \mathbf{p}$$

$$\mathcal{M} \propto \Lambda^5 \int dr u(r) r e^{-\frac{r^2 \Lambda^2}{2}} v(r) \rightarrow \tilde{d}_0(\Lambda) \text{ (for large } \Lambda)$$

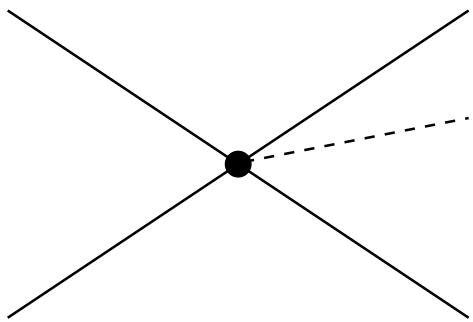


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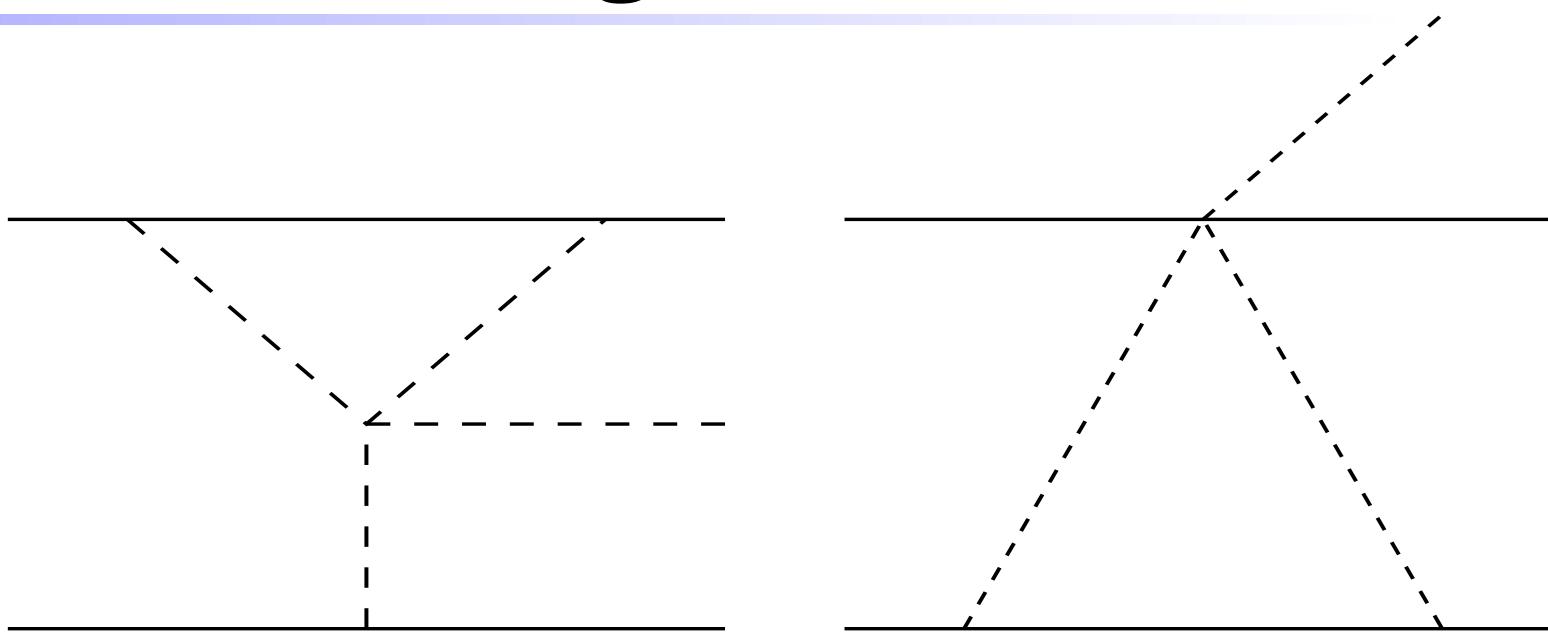


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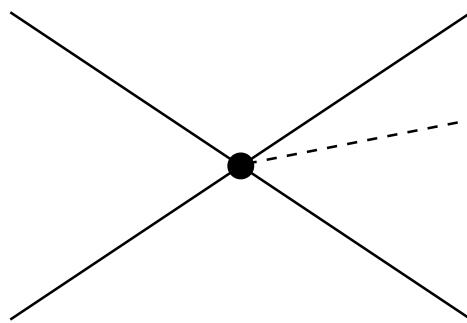
$$\mathcal{M} \propto \Lambda^5 \int dr u(r) r e^{-\frac{r^2 \Lambda^2}{2}} v(r) \rightarrow \tilde{d}_0(\Lambda) \text{ (for large } \Lambda)$$

- mixing DR and cut-off?
- DR/cut-off for matrix elements? $V_{\text{low}k}$ useful?
- Using cut-off+counter term to ren $\mathcal{O} \Rightarrow$ autoren of \mathcal{M} ?

Surviving s -wave NLO



Need counter-term to ‘restore’ χS for \mathcal{M} :



Cf. Andreas's talk on NN P -waves Tuesday this week

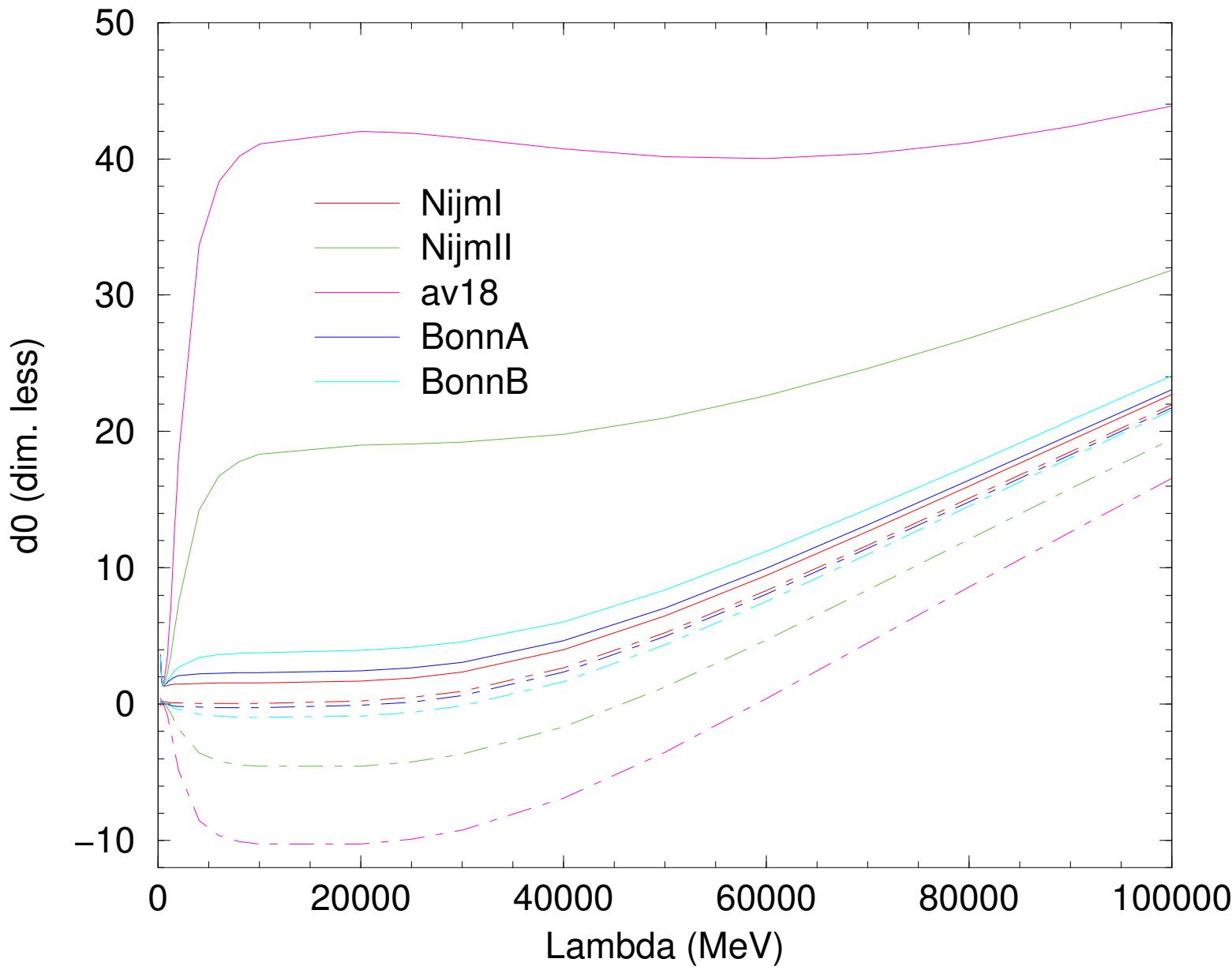
Nogga, Timmermans, and van Kolck, nucl-th/0506005

Results without Δ (prel.)

For each Λ there are 2 solutions for \tilde{d}_0 :

| pot | $\Lambda = 500 \text{ MeV}$ | | $\Lambda = 800 \text{ MeV}$ | | $\Lambda = 1200 \text{ MeV}$ | |
|--------|-----------------------------|---------|-----------------------------|--------|------------------------------|---------|
| NijmI | 1.342 | 0.198 | 1.345 | 0.156 | 1.414 | 0.121 |
| NijmII | 1.480 | 0.0618 | 2.038 | -0.162 | 3.495 | -0.648 |
| av18 | 1.626 | -0.0952 | 2.931 | -0.602 | 6.944 | -1.865 |
| Bonn A | 1.355 | 0.108 | 1.525 | 0.0127 | 1.795 | -0.0826 |
| Bonn B | 1.383 | 0.125 | 1.612 | 0.0122 | 2.053 | -0.1498 |

$\tilde{d}_0(\Lambda)$ at REALLY high Λ



Conclusions and outlook

- Understand CS XS before tackling CSB!
- CS diagrams known to NLO
- Many *s*-wave diagrams, but most loops cancel
- Few *p*-wave diagrams, no NLO!
- No new scale with explicit Δ !
- Loops wf dep. NEED COUNTER TERM!
- Implementation of 1-body term
- Leakage $\alpha \rightarrow \beta$, new experiments?
- Inelasticities, Δ and N^*

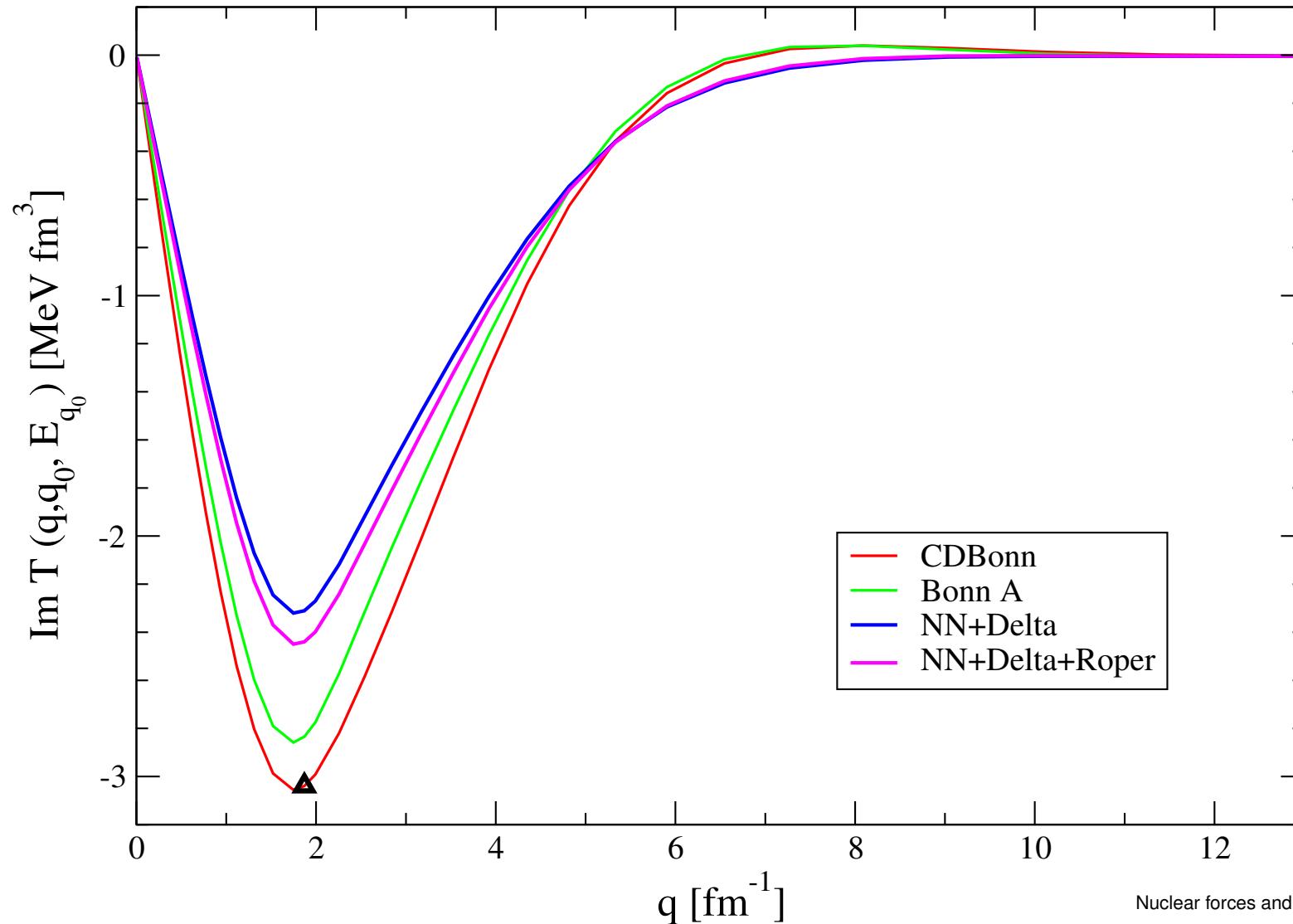
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- Leakage $\alpha \rightarrow \beta$, new experiments?
- Inelasticities, Δ and N^*
- Use chiral wfs, $V_{\text{low}k}$
- Develop CSB *s*- and *p*-waves! Counter terms?
- CSB with explicit Δ ?

Imaginary $T(^3P_1)$

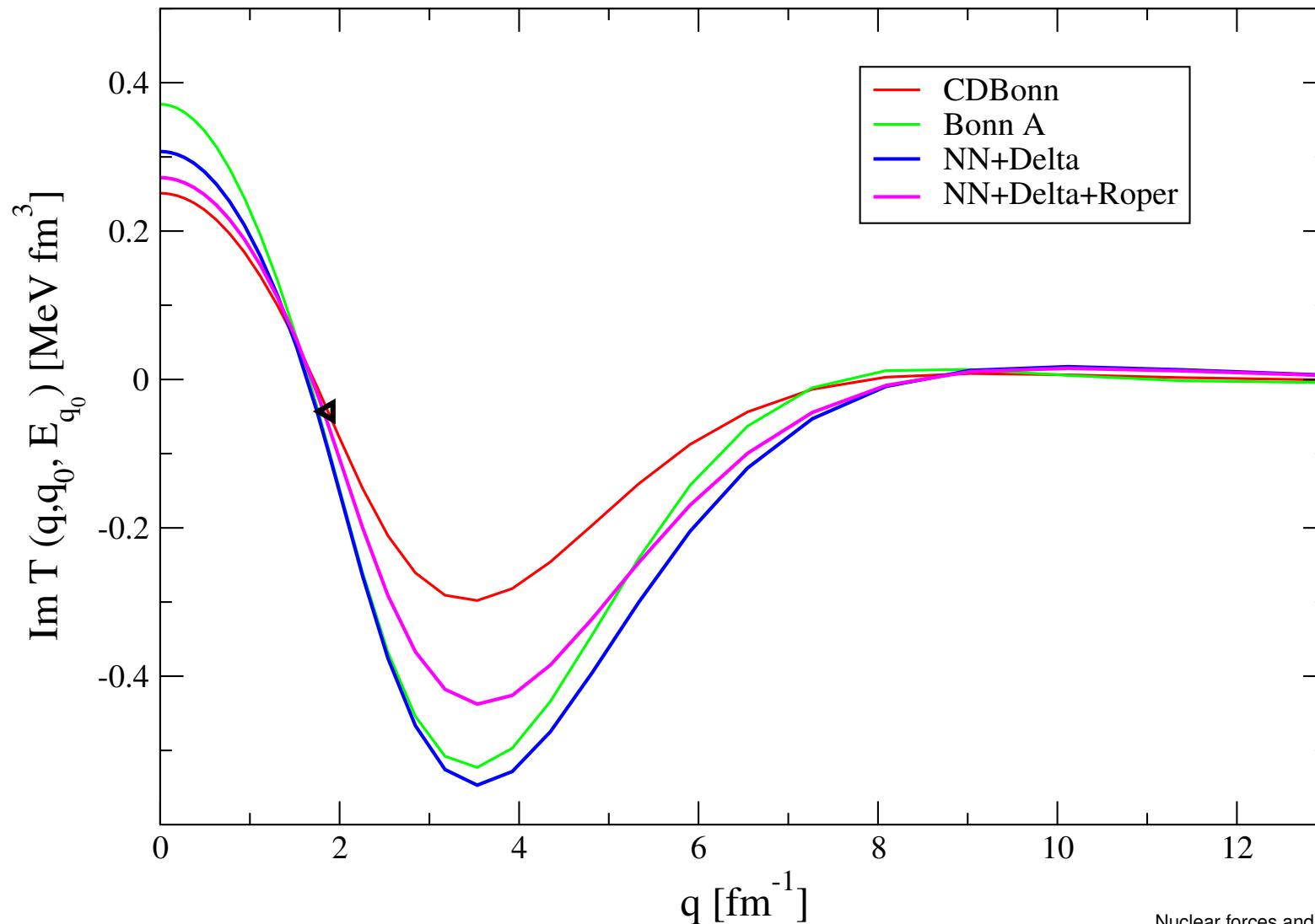


3P_1 half-shell T-matrix ($E_{\text{lab}}=290 \text{ MeV}$)



Imaginary $T(^1S_0)$

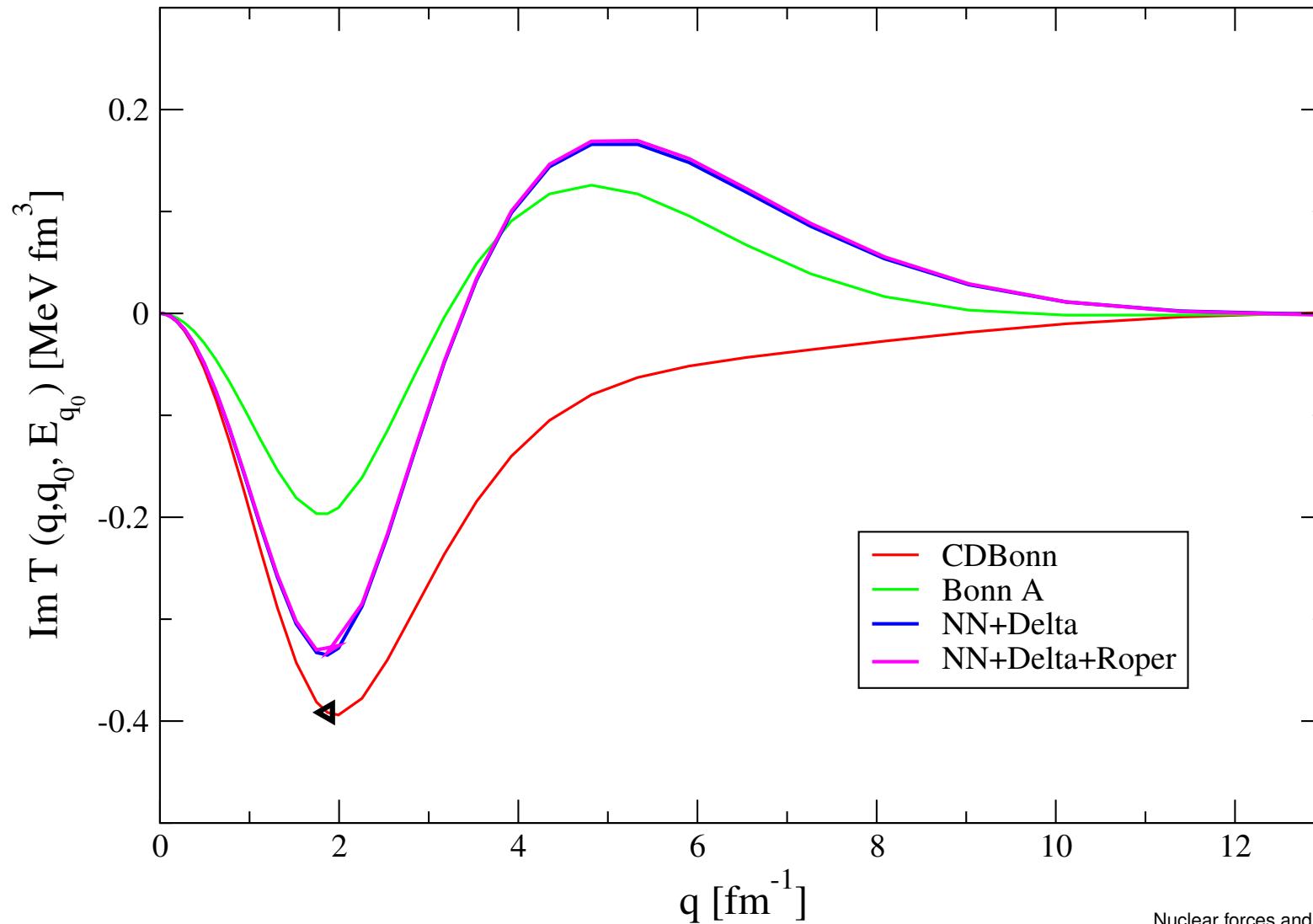
1S_0 half-shell T-matrix ($E_{\text{lab}} = 290 \text{ MeV}$)



Imaginary $T(^1D_2)$



1D_2 half-shell T-matrix ($E_{\text{lab}}=290$ MeV)



Power counting (CSB)

Relevant new scales:

- $(M_\Delta - M)/M$, Δ -nucleon mass difference
- $\epsilon = (m_d - m_u)/(m_d + m_u)$
- α from photon exchanges

Typical CSB parameters:

$$\epsilon \frac{m_\pi^2}{\Lambda_\chi^2} \sim 1\% \text{ — strong: } \delta M, \beta_1, \gamma_1$$

$$\frac{\alpha}{\pi} \sim 0.25\% \text{ — EM: } \bar{\delta} M, \bar{\beta}_3, \bar{\gamma}_3$$

Also soft photon exchanges contribute

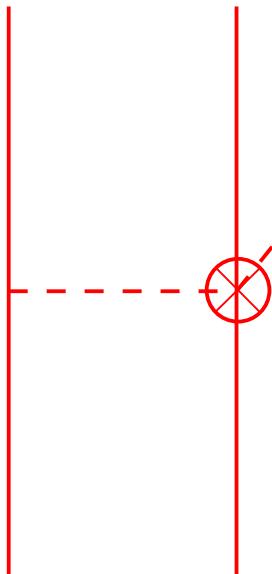
CSB Lagrangian

$$\begin{aligned}\mathcal{L}_{CSB} = & \frac{\delta M}{2} N^\dagger \left(\tau_3 - \frac{\pi_3 \boldsymbol{\tau} \cdot \boldsymbol{\pi}}{2f_\pi^2} \right) N - \frac{\beta_1 + \bar{\beta}_3}{2f_\pi} N^\dagger \vec{\sigma} \cdot \vec{\nabla} \pi_3 N \\ & + \frac{\bar{\delta}M}{2} N^\dagger \left(\tau_3 + \frac{\pi_3 \boldsymbol{\tau} \cdot \boldsymbol{\pi} - \boldsymbol{\pi}^2 \tau_3}{2f_\pi^2} \right) N \\ & - \frac{(\gamma_1 + \bar{\gamma}_3)}{2f_\pi} N^\dagger N \left\{ N^\dagger \vec{\sigma} N \cdot \vec{\nabla} \pi_3 - \frac{1}{2M} \left[i N^\dagger \dot{\pi}_3 \vec{\sigma} \cdot \vec{\nabla} N + \text{h.c.} \right] \right\} \\ & + \dots\end{aligned}$$

Undetermined coefficients!

$$\delta M, \bar{\delta}M, \beta_1, \bar{\beta}_3, \gamma_1, \bar{\gamma}_3$$

CSB *s*-wave LO



$$m_n - m_p = \delta M + \bar{\delta}M = 1.29 \text{ MeV.}$$

Natural sizes:

$$\delta M \sim \epsilon m_\pi^2 / \Lambda_{\text{QCD}} \sim 7 \text{ MeV}$$

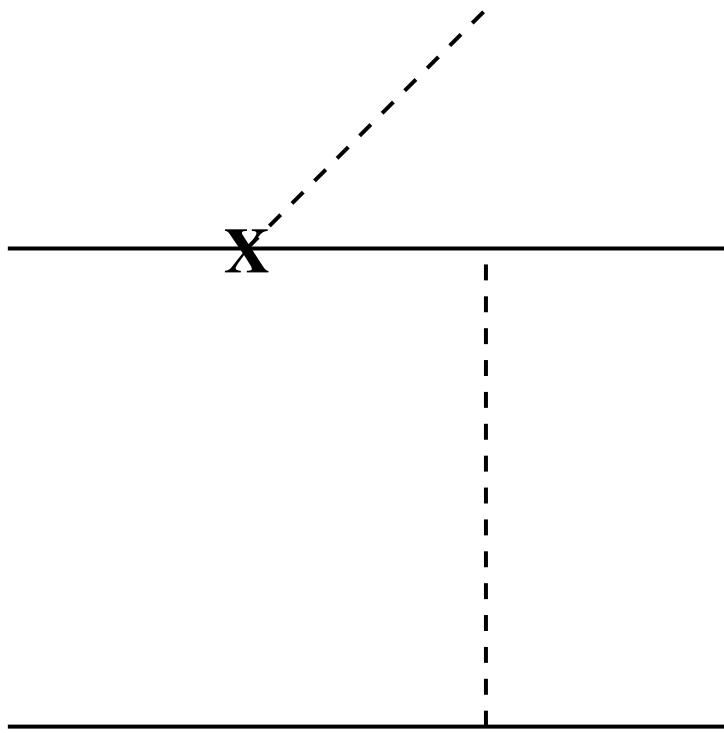
$$\bar{\delta}M \sim -\alpha \Lambda_{\text{QCD}} / \pi \sim -2 \text{ MeV}$$

$$\epsilon = (m_d - m_u) / (m_d + m_u) \sim 1/3$$

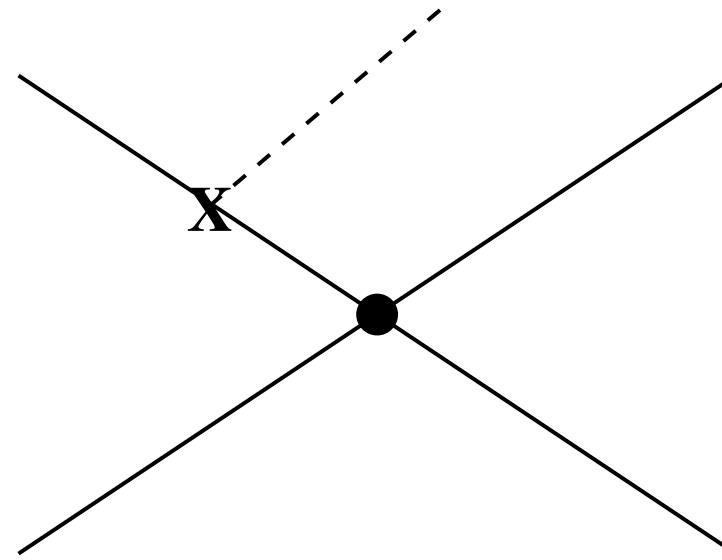
$$\text{CSB parameter } (\delta M - \frac{1}{2}\bar{\delta}M) / 2f_\pi \sim 22 \cdot 10^{-3}$$

Soft photons gives other LO and NLO amplitudes

Short-range interactions (HME) at NNLO



(a)



(b)

Under development!