

1999

There once was a workshop at Trento
to Effectives' and Potentials' memento.

Discussions galore
spilt blood on the floor –
and Bira told me 't was meant to.

Three-Nucleon Forces in Effective Field Theory without Pions

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- 1 A Few Nucleons at Very Low Energies
- 2 The Three-Nucleon System to All Orders
- 3 Implications and Applications
- 4 Concluding Questions

How to root Nuclear Physics in QCD?

How to systematise an EFT with fine-tuning?

Bedaque/hg/Hammer/Rupak: *Nucl. Phys.* **A714** (2003), 589 [nucl-th/0207034]

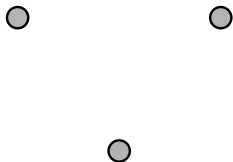
hg: *Nucl. Phys.* **A744** (2004), 192 [nucl-th/0404073]; *Nucl. Phys.* A in press [nucl-th/0502039]

Mathematica code: <http://www.ph.tum.de/~hgrie>

1. A Few Nucleons at Very Low Energies

Pictures of the nucleon:

$$p_{\text{typ.}} \lesssim 100 \text{ MeV} , \Delta x \gtrsim 2 \text{ fm}$$

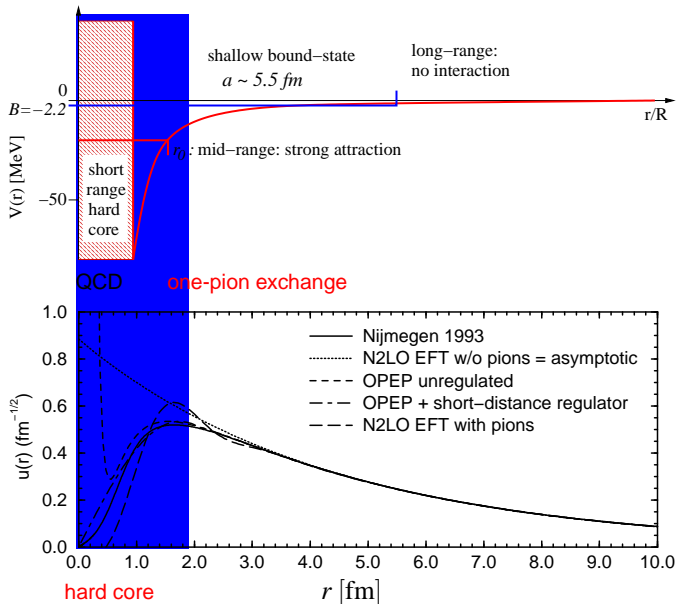


A microscope with very bad resolution!

Non-trivial: Few-nucleon systems with size $a \sim 5 \text{ fm} \gg 1 \text{ fm} \sim R_N$ constituent size!

(a) It's Natural to Have Unnatural Scales

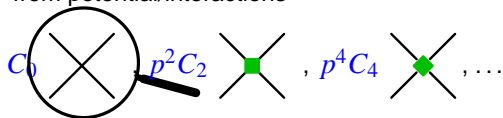
Deuteron bound-state **fine-tuned**: $B = 2.225 \text{ MeV}$, size $a \approx 5 \text{ fm} \gg R \approx 1 \text{ fm}$



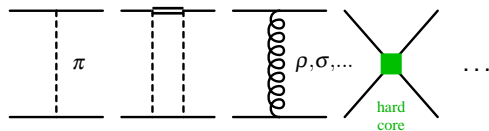
Wave-function at large distances $r \gg R$:

$$\Psi \rightarrow Z \frac{e^{-r/a}}{r} \left[1 + c_1 \frac{R}{r} + c_2 \left(\frac{R}{r} \right)^2 + \dots \right]$$

from potential/interactions



encode QCD: **minimal number** of parameters



$$\Rightarrow \text{systematic expansion in } Q = \frac{\text{target size } R}{\text{resolution } a} \approx \frac{1}{3 \text{ to } 5} \ll 1$$

“EFT without pions” EFT(π)

conceptually well-understood \Rightarrow **predictive power, error estimate**

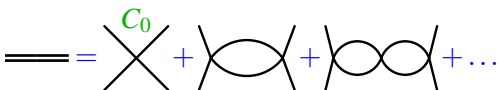
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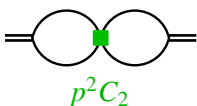
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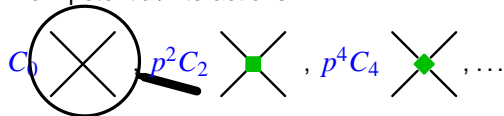
Coefficients C_{2n} from simple observables:

LO (30%): 

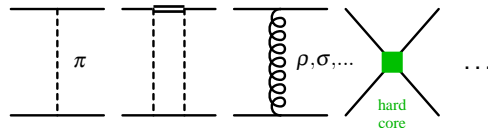
NLO (10%):  effective ranges ρ_0

NNLO (3%): $\rightarrow d$ -wave contributions etc.

from potential/interactions



encode QCD: **minimal number** of parameters



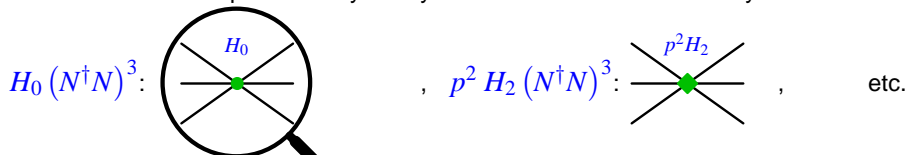
$$\Rightarrow \text{systematic expansion in } Q = \frac{\text{target size } R}{\text{resolution } a} \approx \frac{1}{3 \text{ to } 5} \ll 1$$

“EFT without pions” EFT(π)

conceptually well-understood \Rightarrow **predictive power, error estimate**

(b) Three-Body Forces in EFT(π)

EFT: All interactions permitted by the symmetries of QCD \implies 3-body forces



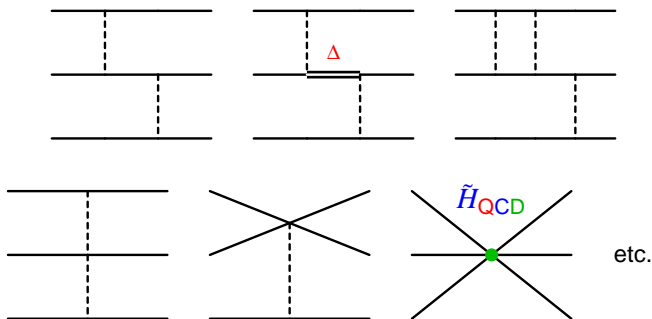
\implies Which channels and observables most sensitive?

How important?



At which order in Q do they start to contribute?

What are they?

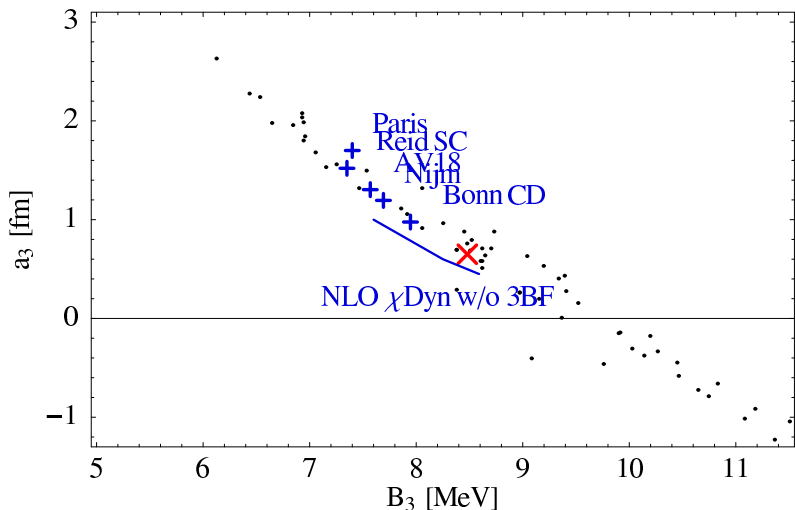


Generic problem: Predictions by modern, high-precision *NN*-potentials differ vastly, but all on Phillips line (1969).

Ad-hoc three-body forces make up for difference.

How to predict 3-body forces?

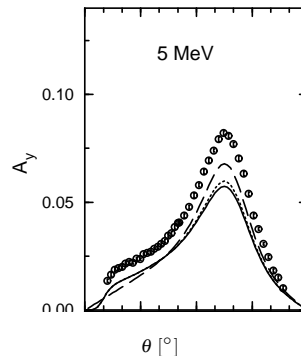
EFT($\not{\pi}$): universal, a priori one free parameter $H_0(\Lambda)$.



×: exp. +: modern *NN* potentials ∴: more potentials

More Problems:

A_y-problem: *Nd* spin-observable



nd → *tγ* at thermal energies:

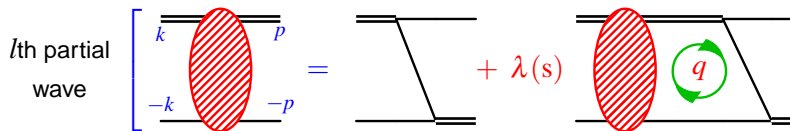
Pot. Mod.'s: $\sigma = [0.49 \dots 0.66]$ mb

exp.: $\sigma = [0.508 \pm 0.015]$ mb

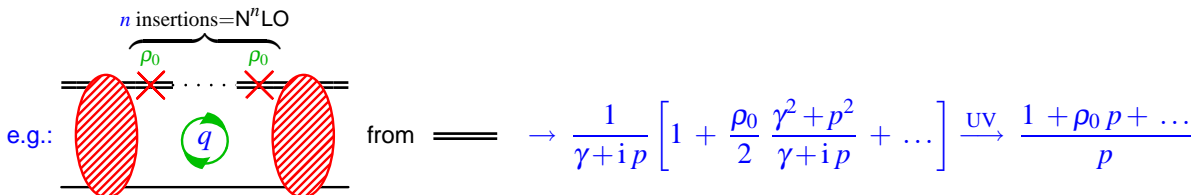
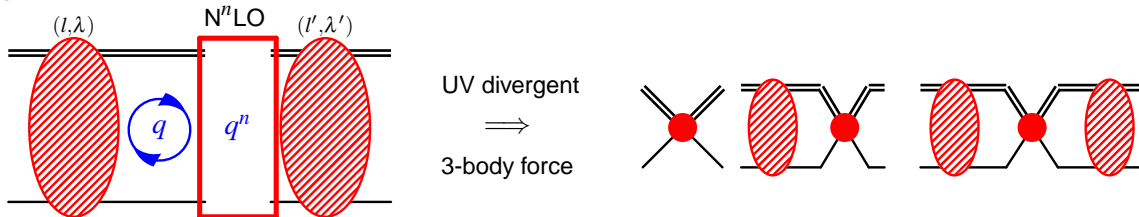
(a) Leading Three-Body Forces from Naïve Dimensional Analysis

Include 3BF if and only if needed as counter-term to cancel cut-off dependence of low-energy observables.

LO: Faddeev integral equation for half off-shell amplitude $\mathcal{A}(k, q)$: (Skorniakov/Ter-Martirosian 1957)



Higher Orders: Perturbation around LO



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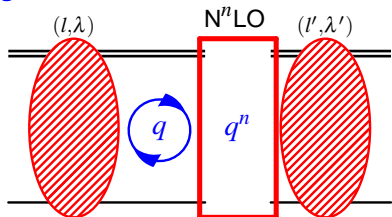
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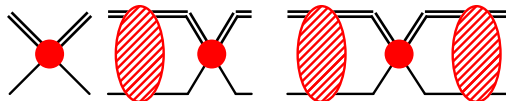
(Skorniakov/Ter-Martirosian 1957)

$$l\text{th partial wave} \left[\begin{array}{c} k \\ -k \end{array} \right] \begin{array}{c} p \\ -p \end{array} = \text{tree} + \lambda(s) \text{loop} \propto \frac{1}{p^{s_l(\lambda)+1}}$$

Higher Orders: Perturbation around LO



UV divergent
 \implies
 3-body force



$$\text{UV} \rightarrow \frac{q^5}{q q^2} \frac{1}{q^{s_l(\lambda)+1}} \frac{1}{q^{s_{l'}(\lambda')+1}} q^n = q^{n - s_l(\lambda) - s_{l'}(\lambda')}$$

3-body force counter-term $\iff n \geq \text{Re}[s_l(\lambda) + s_{l'}(\lambda')]$

3-body force counter-term $\iff n \geq \text{Re}[s_l(\lambda) + s'_l(\lambda')]$

l th partial wave

$$= \text{tree} + \lambda(s) \text{loop} \propto \frac{1}{p^{s_l(\lambda)+1}} ; \text{ simplistic: } \frac{1}{p^2} \frac{k^l}{p^l}$$

Simplistic argument: Amplitude in UV like potential/Born series (e.g. Weinberg's proposal)

Example S-wave:

$$\frac{1}{p^2} + \frac{q^5}{q^{2 \times 3} q} \sim \frac{1}{q^2} + \frac{q^{5 \times 2}}{q^{2 \times 5} q q} \sim \frac{1}{q^2} + \dots$$

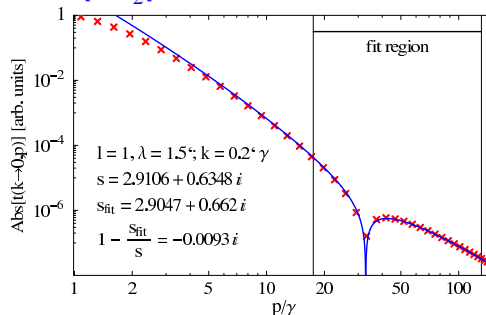
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UV ($k, \gamma \ll p, q$) decoupled by Wigner's SU(4)-symmetry

\implies resolvent $a_{(l,s)}(p) = (-)^l \frac{4 \lambda(s)}{\sqrt{3} \pi} \int_0^\infty \frac{dq}{p} a_{(l,s)}(q) Q_l \left[\frac{p+q}{q} \right]$ by Mellin trafo: $a_{(l,s)}(p) \propto \frac{1}{p^{s_l(\lambda)+1}}$

$$1 = (-)^l \frac{2^{2-l} \lambda}{\sqrt{3} \pi} \frac{\Gamma[\frac{l+1+s_l(\lambda)}{2}] \Gamma[\frac{l+1-s_l(\lambda)}{2}]}{\Gamma[l+\frac{3}{2}]} {}_2F_1\left[\frac{l+1+s_l(\lambda)}{2}, \frac{l+1-s_l(\lambda)}{2}; l+\frac{3}{2}; \frac{1}{4}\right]$$



3-body force counter-term $\iff n \geq \text{Re}[s_l(\lambda) + s'_l(\lambda')]$

l th partial wave $\left[\begin{array}{c} \text{---} k \text{---} p \\ \text{---} -k \text{---} -p \end{array} \right] = \text{---} \text{---} + \lambda(s) \left[\begin{array}{c} \text{---} \text{---} \\ \text{---} \text{---} \end{array} \right] \propto \frac{1}{p^{s_l(\lambda)+1}} ; \text{ simplistic: } \frac{1}{p^2} \frac{k^l}{p^l}$

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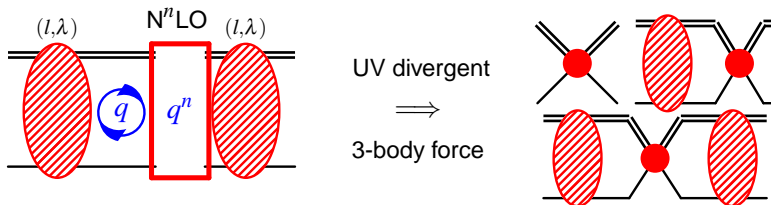
$$1 = (-)^l \frac{2^{1-l} \lambda}{\sqrt{3}\pi} \frac{\Gamma[\frac{l+1+s_l(\lambda)}{2}] \Gamma[\frac{l+1-s_l(\lambda)}{2}]}{\Gamma[l+\frac{3}{2}]} {}_2F_1\left[\frac{l+1+s_l(\lambda)}{2}, \frac{l+1-s_l(\lambda)}{2}; l+\frac{3}{2}; \frac{1}{4}\right]$$

	$s_l(\lambda)$	$l=0$	$l=1$	$l=2$	$l=3$
simplistic: $s_l = l+1$		1	2	3	4
Quartet/Wigner-antisym. Doublet ($s = \frac{3}{2}$): $\lambda = -\frac{1}{2}$		2.16	1.77	3.10	4.04
Wigner-sym. Doublet ($s = \frac{1}{2}$): $\lambda = 1$		$\pm 1.0062 i$	2.86	2.82	3.92

\implies UV-scaling not as guessed.

(c) Leading Three-Body Forces in the 3-Nucleon Channels

Include 3BF if and only if needed as counter-term to cancel cut-off dependence of low-energy observables.



channel partial waves	naïve dim. an. $\text{Re}[s_l(\lambda) + s_{l'}(\lambda')]$	simplistic $l + l' + 2$		typ. size $1/3^n$
${}^2S_{Ws} - {}^2S_{Ws}$	LO	$N^2\text{LO}$	promoted	100%
${}^2S_{Ws} - {}^2S_{Wa}$	$N^{2.2+2}\text{LO}$	$N^{2+2}\text{LO}$	demoted	1%
${}^2S_{Wa} - {}^2S_{Wa}$	$N^{4.3+2}\text{LO}$			0.1%
${}^2S_{Ws} - {}^4D$	$N^{3.1}\text{LO}$	$N^4\text{LO}$	promoted	3%
${}^2S_{Wa} - {}^4D$	$N^{5.3}\text{LO}$		demoted	0.3%
${}^2P_{Ws} - {}^2P_{Ws}$	$N^{5.7}\text{LO}$	$N^4\text{LO}$	demoted	0.2%
${}^2P_{Ws} - {}^2P_{Wa}, {}^2P_{Ws} - {}^4P$	$N^{4.6}\text{LO}$		demoted	0.6%
${}^2P_{Wa} - {}^2P_{Wa}, {}^4P - {}^4P$	$N^{3.5}\text{LO}$			2%
${}^4S - {}^4S$	$N^{4.3+2}\text{LO}$	$N^{2+2}\text{LO}$	demoted	0.1%
${}^4S - {}^2D_{Ws}$	$N^{5.0}\text{LO}$	$N^4\text{LO}$	demoted	0.4%
${}^4S - {}^2D_{Wa}, {}^4S - {}^4D$	$N^{5.3}\text{LO}$		demoted	0.3%
higher	\sim as simplistic	$N^{l+l'+2}\text{LO}$		

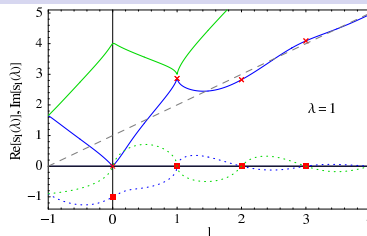
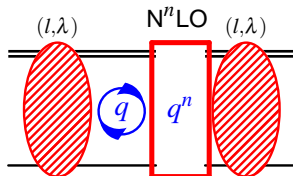
The Three Big Lies of Nuclear Physics

Nuclear Power is Safe.

They have Weapons of Mass Distruction.

My Power-Counting is Systematic.

(d) Context and Caveats



Context

- Extended **Naïve Dimensional Analysis** to **non-perturbative** field theories.
- **Criterion**: Independence of short-distance details at each order.
- Only phenomenological input: **shallow 2-body bound states** \implies LO non-perturbative.
- **Higher partial waves** close to perturbative estimate.
- **Higher orders**: 3body-forces with $2m$ more derivatives at $n \geq \text{Re}[s_l(\lambda) + s'_l(\lambda')] - 2m$;
- **External currents** straightforward: replace kernel.

Renormalisability.

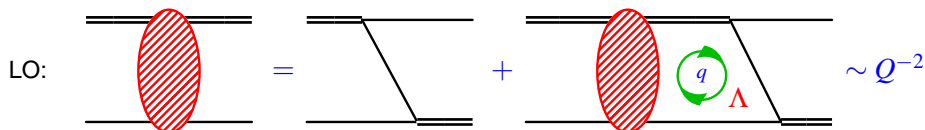
Caveats

- **No more re-summations** after LO No contrary cases known.
- **3-body force** \iff **counter-term**: minimal number. Assumed natural size set by running with cut-off.
- Only **superficial degree of divergence**. \implies Future: Construct 3body forces!
- Cancellations between diagrams? Overlapping divergences? \implies **Demote 3body forces.**
- Generically **fractional orders**: \implies Clear criterion?
- ${}^4S_{\frac{3}{2}}$: $s = 2.16 \implies$ 3BF at $N^{6.3}$ LO Diverges at N^7 LO; converges at N^6 LO very slowly.
- 4F : $s = 5.02 \implies$ 3BF at $N^{10.04}$ LO Diverges at N^{11} LO; converges at N^{10} LO very slowly.

That's nice, dear.

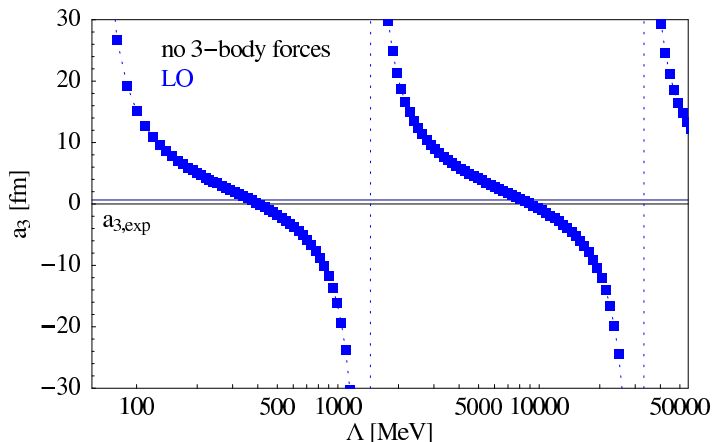
3. Implications and Applications

(a) The Problem: nd -Scattering, ${}^2S_{\frac{1}{2}}$ Wave (“triton channel”)



\downarrow \rightarrow $\uparrow\uparrow$: No Pauli principle, no centrifugal barrier \Rightarrow 3-body forces at N2LO?

$H_0 \sim Q^0$



Slight cut-off variation has dramatic effect on low-energy observables.

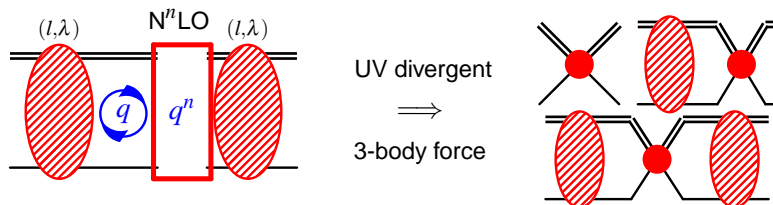
Danilov, Minlos/Faddeev 1961

$$s_0({}^2S_{\frac{1}{2}}) = \pm 1.006 \dots i \Rightarrow$$

On-shell depends on UV-phase δ :

$$A_{2S_{\frac{1}{2}}}(k \rightarrow 0, p) \propto \frac{\cos[1.0062 \ln p + \delta]}{p}$$

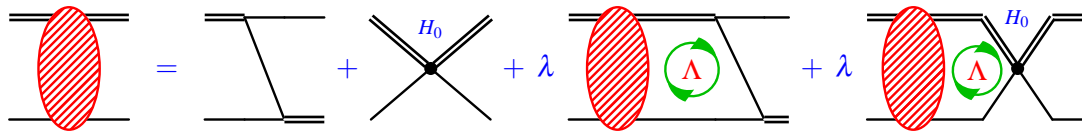
Include 3BF if and only if needed as counter-term to cancel cut-off dependence of low-energy observables.



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${}^2S_{Ws} - {}^2S_{Ws}$ ${}^2S_{Ws} - {}^2S_{Wa}$ ${}^2S_{Wa} - {}^2S_{Wa}$	LO $N^{2,2+2}LO$ $N^{4,3+2}LO$	N^2LO $N^{2+2}LO$	promoted demoted	100% 1% 0.1%
${}^2S_{Ws} - {}^4D$ ${}^2S_{Wa} - {}^4D$	$N^{3,1}LO$ $N^{5,3}LO$	N^4LO	promoted demoted	3% 0.3%
${}^2P_{Ws} - {}^2P_{Ws}$ ${}^2P_{Ws} - {}^2P_{Wa}$, ${}^2P_{Ws} - {}^4P$ ${}^2P_{Wa} - {}^2P_{Wa}$, ${}^4P - {}^4P$	$N^{5,7}LO$ $N^{4,6}LO$ $N^{3,5}LO$	N^4LO	demoted demoted	0.2% 0.6% 2%
${}^4S - {}^4S$ ${}^4S - {}^2D_{Ws}$ ${}^4S - {}^2D_{Wa}$, ${}^4S - {}^4D$	$N^{4,3+2}LO$ $N^{5,0}LO$ $N^{5,3}LO$	$N^{2+2}LO$ N^4LO	demoted demoted demoted	0.1% 0.4% 0.3%
higher	~ as simplistic	$N^{l+l'+2}LO$		

Tenet: Include specific 3BF if and only if needed to cancel off-shell dependence of observables.

Momentum-independent **three-body** force $H_0(\Lambda)$ must be LO to absorb cut-off dependence.

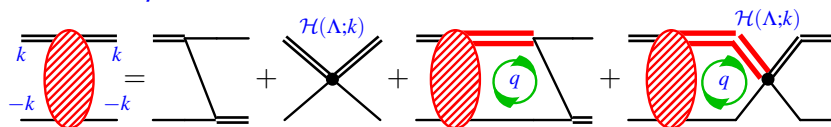


Tune such that \mathcal{A} cut-off independent in UV.

Include specific 3-body force $H_{2n}(\Lambda)$ **if and only if** needed as **counter-term** to cancel cut-off dependence of low-energy observables.

UV limit:

$k \sim \gamma \ll \frac{1}{\rho_0} \ll \Lambda, q$ (off-shell mom.) \implies Lowest 3-body forces **Wigner-SU(4)-symmetric**.



LO

NLO

N2LO

$$\Psi_d(q \rightarrow \infty) \equiv \rightarrow \sqrt{\frac{4}{3}} \frac{1}{q} + \frac{4\gamma}{3q^2} + \frac{\rho_0}{2} + \frac{3k^2 + 4\gamma^2}{3\sqrt{3}q^3} + \frac{2\gamma\rho_0}{\sqrt{3}q} + \frac{\sqrt{3}}{8} q\rho_0^2$$

$$V_{Nd}(q \rightarrow \infty) \equiv \rightarrow \frac{1}{q^2} + \frac{k^2 - 12\gamma^2}{12q^4}$$

$$\implies \mathcal{H} \equiv H_0^{\text{LO}}(\Lambda) + H_0^{\text{NLO}}(\Lambda) + \frac{k^2}{\Lambda^2} H_2(\Lambda) + H_0^{\text{N2LO}}(\Lambda)$$

LO and NLO (< 10% accuracy):

One free parameter H_0 ,
fixed e.g. by triton binding energy.

N2LO and N3LO (< 1% accuracy):

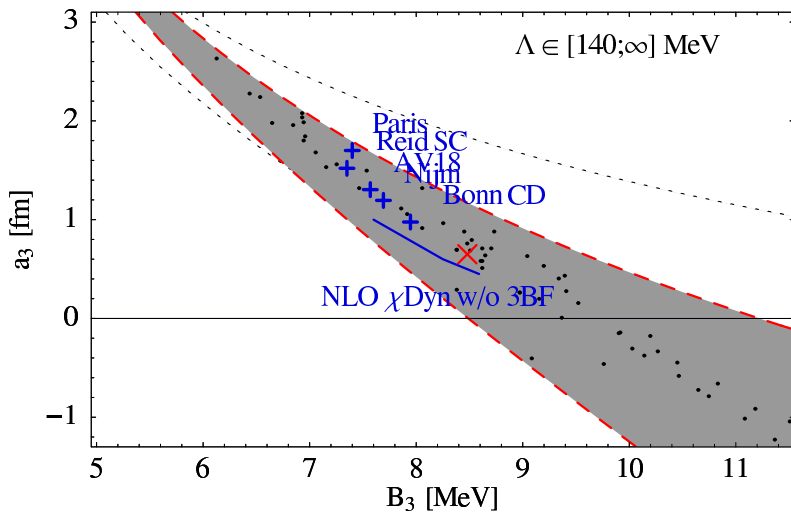
One more free parameter H_2 ,
fixed best by **scattering length**.

Generic problem: Predictions by modern, high-precision NN -potentials differ vastly, but all on Phillips line (1969).

Ad-hoc three-body forces make up for difference.

How to predict 3-body forces?

EFT(π): universal, a priori one free parameter $H_0(\Lambda)$.



×: exp.

+ : modern NN potentials

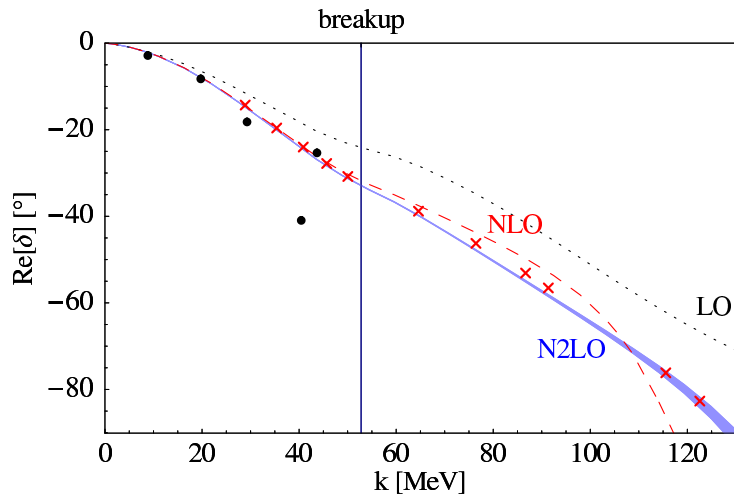
· : more potentials

(c) Doublet-S Wave nd Phase Shift

×: AV18+U IX (Kievski 2002)

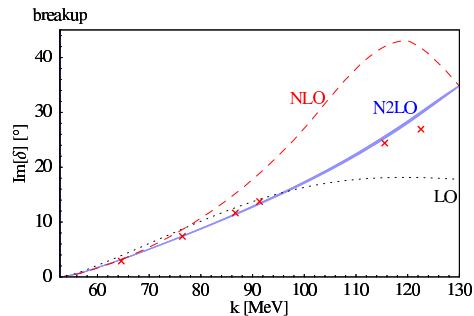
●: PWA 1967 (Seagrave/van Oers)

■: N2LO, $\Lambda \in [200; \infty]$ MeV



Fix H_0 to one observable (a_3).

N2LO: H_0 & H_2 to (B_3 , a_3)



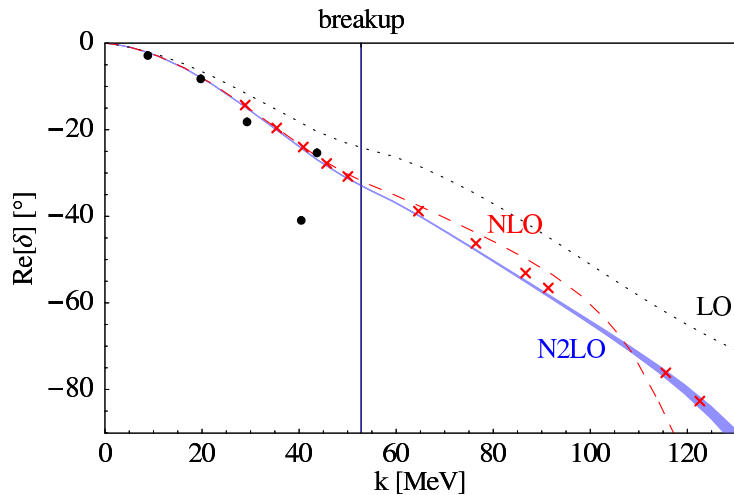
Agrees well with sophisticated, modern potential model calculations.

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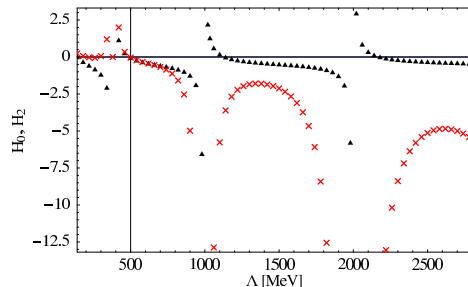
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Variation of 3body force



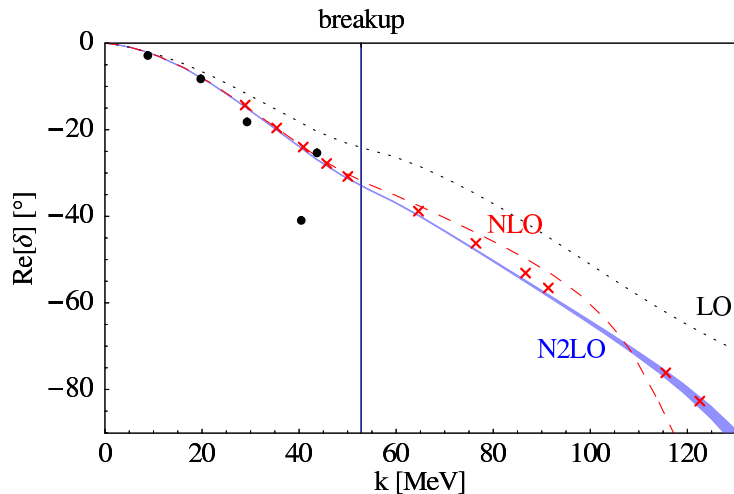
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Agrees well with sophisticated, modern potential model calculations.

- Convergence to **Nature**.
- Order by order smaller **corrections**.
- Order by order smaller **cut-off dependence**.

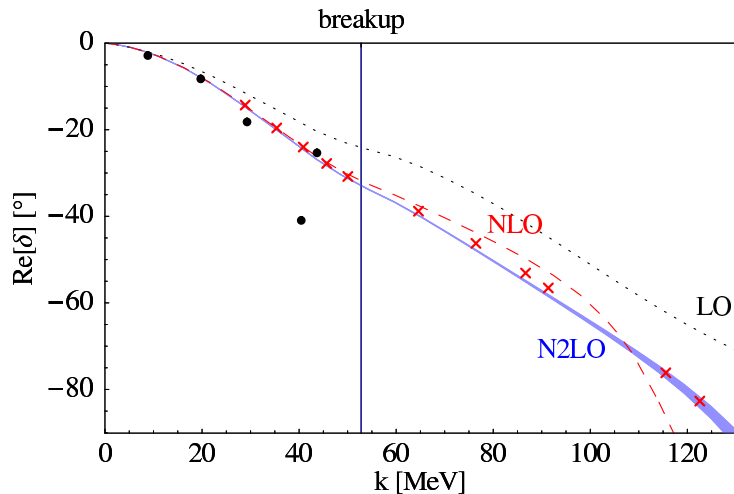
No Nature here.

O.k.

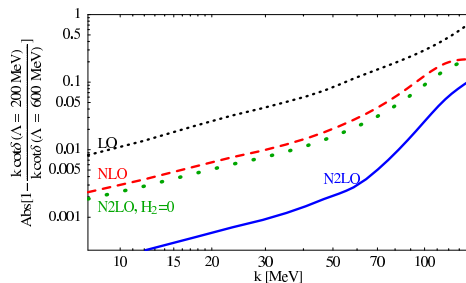
“Lepage plot”.

×: AV18+U IX (Kievski 2002)

●: PWA 1967 (Seagrave/van Oers)

N2LO: $\Lambda \in [200; \infty]$ MeV


“Lepage-plot”: relative, log. cut-off dep.



Agrees well with sophisticated, modern potential model calculations.

$$\left| 1 - \frac{k \cot \delta(\Lambda = 200 \text{ MeV})}{k \cot \delta(\Lambda = \infty)} \right| \sim \underbrace{\left(\frac{p_{\text{typ.}}}{\Lambda_{\text{fit}}} \right)^n}_{Q^n}$$

 \Rightarrow Fit to $k \in [70; 100 \dots 130]$ MeV

	LO	NLO	N ² LO	N ² LO without H_2
n fitted	~ 1.9	2.9	4.8	3.1
n expected	2	3	4	4!!

(d) Simpler Picture: The Three-Nucleon System in Coordinate Space

Fourier:
$$\left[-\frac{1}{R} \frac{\partial}{\partial R} R \frac{\partial}{\partial R} + \# \frac{l(l+1) + s_l^2}{R^2} - ME \right] F(R) = 0.$$

$s_l(\lambda)$ **imaginary**, $|s_l(\lambda)| \gg l$ (e.g. $s_0(^2S_{\frac{1}{2}}) = \pm 1.006 \dots i$)

⇒ attractive potential, infinitely many, **deeply** bound states.

⇒ **Thomas & Efimov Effects** (1935, 1971)

⇒ **too much** sensitivity to short-distance

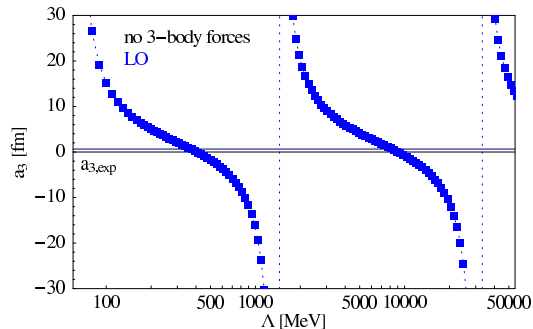
⇒ 3body forces **promoted**

$s_l(\lambda) \gg l+1$ (e.g. $s_0(^4S_{\frac{3}{2}}) = 2.16 \dots \gg 1$)

⇒ more repulsive than estimated

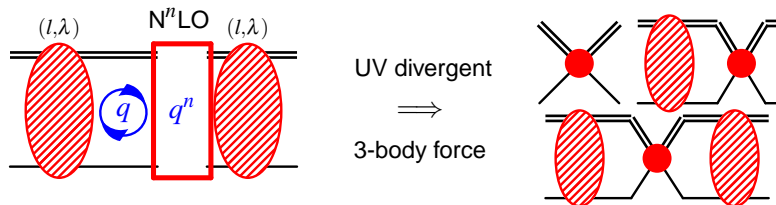
⇒ **less** sensitivity to short-distance

⇒ 3body forces **demoted**



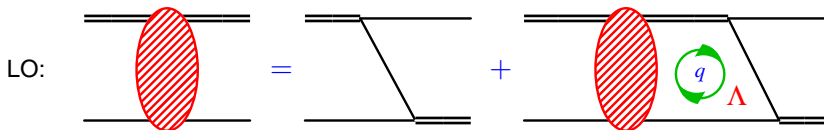
(d) Simpler Picture: The Three-Nucleon System in Coordinate Space

Include 3BF if and only if needed as counter-term to cancel cut-off dependence of low-energy observables.



channel partial waves	naïve dim. an. $\text{Re}[s_l(\lambda) + s_{l'}(\lambda')]$	simplistic $l + l' + 2$		typ. size $1/3^n$
${}^2S_{Ws} - {}^2S_{Ws}$	LO	$N^2\text{LO}$	promoted	100%
${}^2S_{Ws} - {}^2S_{Wa}$	$N^{2.2+2}\text{LO}$	$N^{2+2}\text{LO}$	demoted	1%
${}^2S_{Wa} - {}^2S_{Wa}$	$N^{4.3+2}\text{LO}$			0.1%
${}^2S_{Ws} - {}^4D$	$N^{3.1}\text{LO}$	$N^4\text{LO}$	promoted	3%
${}^2S_{Wa} - {}^4D$	$N^{5.3}\text{LO}$		demoted	0.3%
${}^2P_{Ws} - {}^2P_{Ws}$	$N^{5.7}\text{LO}$	$N^4\text{LO}$	demoted	0.2%
${}^2P_{Ws} - {}^2P_{Wa}, {}^2P_{Ws} - {}^4P$	$N^{4.6}\text{LO}$		demoted	0.6%
${}^2P_{Wa} - {}^2P_{Wa}, {}^4P - {}^4P$	$N^{3.5}\text{LO}$			2%
${}^4S - {}^4S$	$N^{4.3+2}\text{LO}$	$N^{2+2}\text{LO}$	demoted	0.1%
${}^4S - {}^2D_{Ws}$	$N^{5.0}\text{LO}$	$N^4\text{LO}$	demoted	0.4%
${}^4S - {}^2D_{Wa}, {}^4S - {}^4D$	$N^{5.3}\text{LO}$		demoted	0.3%
higher	\sim as simplistic	$N^{l+l'+2}\text{LO}$		

(e) nd scattering length, Quartet-S Wave



N^2 LO EFT($\not{\Lambda}$): $a = [6.35 \pm 0.02] \text{ fm}$

experiment: $[6.35 \pm 0.02] \text{ fm}$

predict first 3BF at $N^{\sim 6}$ LO

$$\approx \pm \left(\frac{1}{3}\right)^3 \times 0.02 \text{ fm} = \pm 0.001 \text{ fm}$$

simplistic: N^4 LO

$$\approx \pm \left(\frac{1}{3}\right)^1 \times 0.02 \text{ fm} = \pm 0.010 \text{ fm}$$

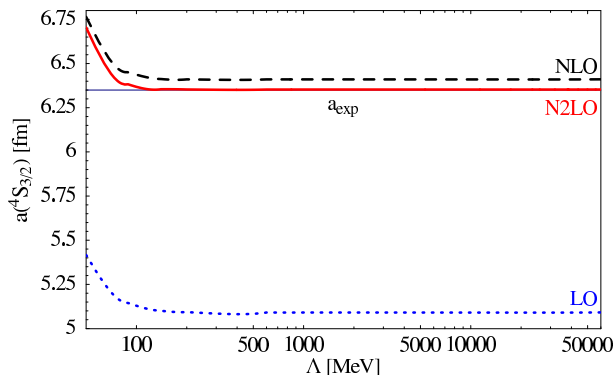
pot. models: $[6.345 \pm 0.002] \text{ fm}$

\Rightarrow input for better $a(^2S)$ from

bound coherent scatt. length: Black et al 2003

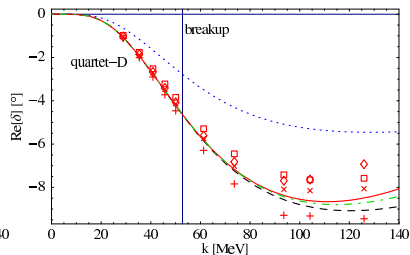
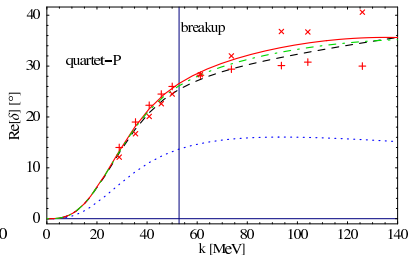
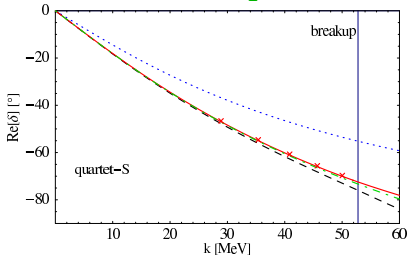
$$a(^2S) = [0.645 \pm 0.003 \pm 0.007_{\text{th}}] \text{ fm}$$

(LO: Skorniakov/Ter-Martirosian 1957, NLO: Efimov 1991, N2LO: Bedaque/van Kolck 1998, hg 2005)

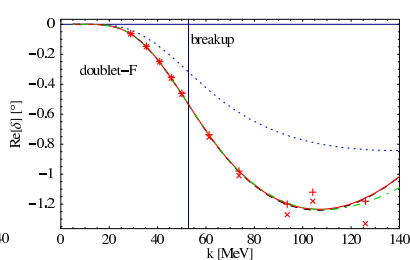
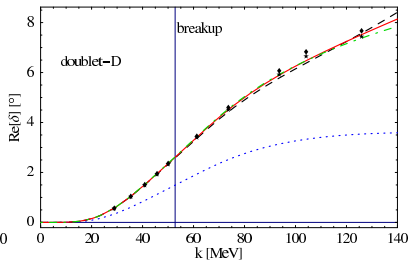
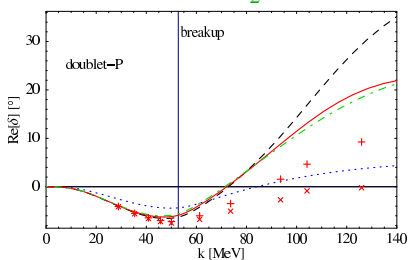


\Rightarrow Alleviates hunt for 3BFs in observables.

Quartet Channel ($s = \frac{3}{2}$)



Doublet Channel ($s = \frac{1}{2}$)




Numerically simple: N2LO code runs within a minute on PC

Agrees well with sophisticated, modern potential model calculations.

N3LO (3-body force!): Splitting/mixing of partial waves $\implies A_y$

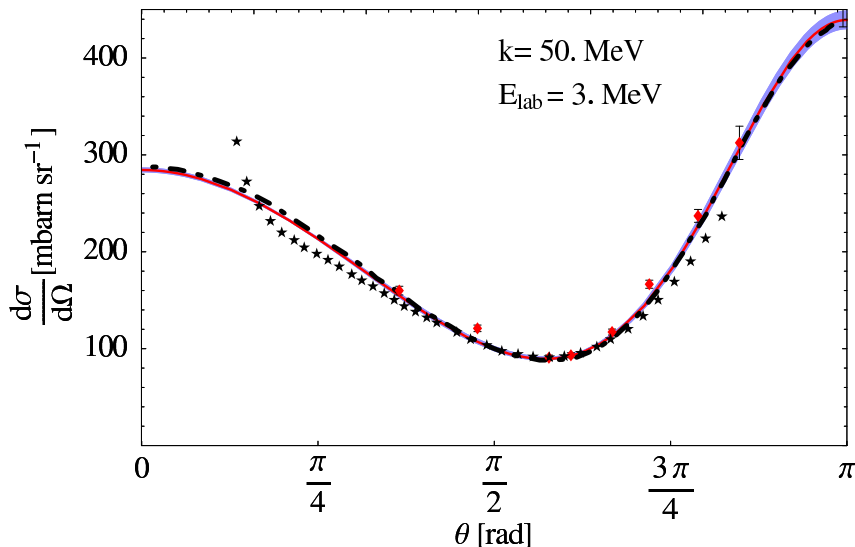
Three-body forces in all other channels even **more suppressed**
than **naively estimated** from centrifugal barrier and Pauli principle.

EFT at N2LO with error bars: 

nd-data: 

pd-data: 

potential models: - - - -



Numerically simple: N2LO code runs within 5 minutes on PC.

Agrees well with sophisticated, modern potential model calculations.

A Problem Solved: $nd \rightarrow t\gamma$ at thermal energies

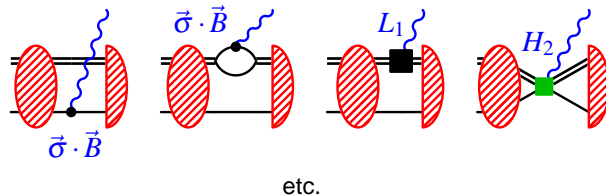
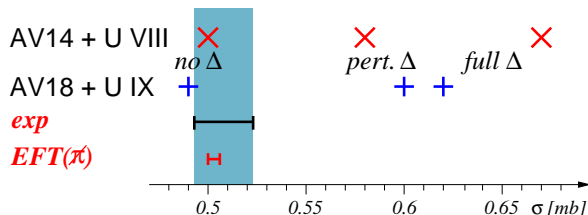
AV 14 + Urbana VIII	no $\Delta(1232)$	0.50 mb	} ?
AV 14 + Urbana VIII	+ $\Delta(1232)$	0.58 mb	
AV 14 + Urbana VIII	perturbative Δ	0.66 mb	
AV 18 + Urbana IX	no $\Delta(1232)$	0.49 mb	
AV 18 + Urbana IX	+ $\Delta(1232)$	0.60 mb	
AV 18 + Urbana IX	perturbative Δ	0.62 mb	

Kievsky/Schiavilla/Viviani 1996, 2004

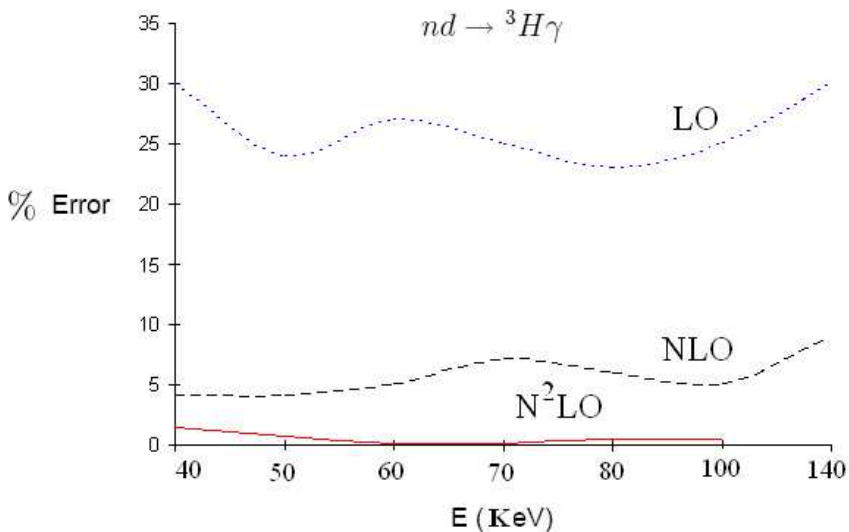
experiment $[0.508 \pm 0.015]$ mb

N^2 LO EFT(π) no new 3BF $[0.503 \pm 0.003]$ mb = $[0.485 + 0.013 + 0.007]$ mb

LO
NLO
N²LO

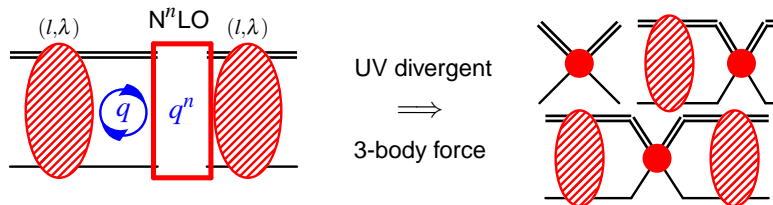


Factor 200 smaller cross-sections.



(i) $nd \rightarrow t\gamma$ for Big-Bang Nucleo-Synthesis

Include 3BF if and only if needed as counter-term to cancel cut-off dependence of low-energy observables.



channel partial waves	naïve dim. an. $\text{Re}[s_l(\lambda) + s_{l'}(\lambda')]$	simplistic $l + l' + 2$		typ. size $1/3^n$
${}^2S_{Ws} - {}^2S_{Ws}$	LO	$N^2\text{LO}$	promoted	100%
${}^2S_{Ws} - {}^2S_{Wa}$	$N^{2,2+2}\text{LO}$	$N^{2+2}\text{LO}$	demoted	1%
${}^2S_{Wa} - {}^2S_{Wa}$	$N^{4,3+2}\text{LO}$		demoted	0.1%
${}^2S_{Ws} - {}^4D$	$N^{3,1}\text{LO}$	$N^4\text{LO}$	promoted	3%
${}^2S_{Wa} - {}^4D$	$N^{5,3}\text{LO}$		demoted	0.3%
${}^2P_{Ws} - {}^2P_{Ws}$	$N^{5,7}\text{LO}$		demoted	0.2%
${}^2P_{Ws} - {}^2P_{Wa}, {}^2P_{Ws} - {}^4P$	$N^{4,6}\text{LO}$	$N^4\text{LO}$	demoted	0.6%
${}^2P_{Wa} - {}^2P_{Wa}, {}^4P - {}^4P$	$N^{3,5}\text{LO}$		demoted	2%
${}^4S - {}^4S$	$N^{4,3+2}\text{LO}$	$N^{2+2}\text{LO}$	demoted	0.1%
${}^4S - {}^2D_{Ws}$	$N^{5,0}\text{LO}$	$N^4\text{LO}$	demoted	0.4%
${}^4S - {}^2D_{Wa}, {}^4S - {}^4D$	$N^{5,3}\text{LO}$		demoted	0.3%
higher	\sim as simplistic	$N^{l+l'+2}\text{LO}$		

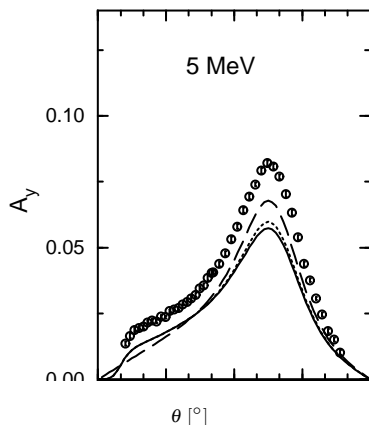
(j) The Future in Few-Nucleon Systems

Classified three-body forces in EFT($\not\hbar$): counter-terms to cancel cut-off dependence.

hg et al.2000-05

An Unsolved Puzzle

A_y -problem: Nd spin-observable



30% discrepancy between all potential models and experiment
 \Rightarrow EFT($\not\hbar$) has to deliver.

The Future:

- ${}^3\text{H}$ and ${}^{3,4}\text{He}$, light nuclei \Rightarrow exotic nuclei, hyper-nuclei
- nuclear & neutrino **Astro-physics**:
big-bang nucleo-synthesis, stellar evolution, Standard Model,...
- **Universality**: molecular systems, Bose-Einstein Condensates,...
- **Partial-wave regularisation?**: answer for any l .
- **Apply methodology to χ EFT**: cf. Nogga/Timmermans/van Kolck 2005
 - LO non-perturbative.
 - Higher orders perturbative.
 - Only** those 2-, 3-, N -body forces to cancel **cut-off dependence**.
 - Check by convergence**: Nature, orders, "Lepage plots".

CTs **promoted or demoted** against Weinberg's proposal.

4. Concluding Questions

EFT(π) is *the* EFT of **QCD** at very low energies: local interactions between nucleons only.

$$\Delta x \gtrsim 2.5 \text{ fm}$$

Systematic classification of all three-body forces:

Tenet: 3BF only as counter-term to cancel cut-off dependence in observables.

3BFs mostly **suppressed or enhanced** against simplistic estimate.

Simplistic power-counting is simplistic.

⇒ **Minimal number of free parameters** at given accuracy.

⇒ **Model-independent, systematic, simple, fast, universal, error-estimates.**

Successful extension of **Naïve Dimensional Analysis** to non-perturbative EFTs.

Plethora of **pivotal physical processes** for **prediction & extraction** of fundamental nucleon properties, e.g.:

$nd \rightarrow t\gamma$, **big-bang nucleosynthesis**, stellar evolution ($E_{\text{typ.}} = 30 - 300 \text{ keV}$)

with Sadeghi

properties of ^3H and ^3He ; **neutrinos and light nuclei** (calibrating SNO)

A_y -problem, hypernuclei, radioactive beams, atomic trimers, BEC, etc.

⇒ **Well on the way to a description of Nuclear Physics deeply rooted in QCD.**